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# The Elusive Present: Hidden Past and Future Dependency and Why We Build Models

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Modeling a temporal process as if it is Markovian assumes the present encodes all of a process's history. When this occurs, the present captures all of the dependency between past and future. We recently showed that if one randomly samples in the space of structured processes, this is almost never the case. So, how does the Markov failure come about? That is, how do individual measurements fail to encode the past? And, how many are needed to capture dependencies between the past and future? Here, we investigate how much information can be shared between the past and future, but not be reflected in the present. We quantify this elusive information, give explicit calculational methods, and draw out the consequences. The most important of which is that when the present hides past-future correlation or dependency we must move beyond sequence-based statistics and build state-based models.

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## I. INTRODUCTION

Until the turn of the nineteenth century, temporal processes were almost exclusively considered to be independently sampled at each time from the same statistical distribution; each sample uncorrelated with its predecessor. These studies were initiated by Jacob Bernoulli in the 1700s [1] and refined by Simeon Poisson [2] and Pafnuty Chebyshev [3] in the 1800s, leading to the weak Law of Large Numbers and the Central Limit Theorem. These powerful results were the first hints at universal laws in stochastic processes, but they applied only to independent, identically distributed (IID) processes—unstructured processes with no temporal correlation, no memory. Moreover, until the turn of the century it was believed that these laws required independence. It fell to Andrei Andreevich Markov (1856–1922) to realize that independence is not necessary. To show this he introduced a new kind of sequence or "chain" of dependent random variables, along with the concepts of transition probabilities, irreducibility, and stationarity [4, 5].

Introducing his "complex chains" in 1907, Markov initiated the modern study of structured, interdependent, and correlated processes. Indeed, in the first and now-famous application of complex chains, he analyzed the pair distribution (2-grams) in the $20,000$ vowels and consonants in Pushkin's poem *Eugeny Onegin* and the $100,000$ letters in Aksakov's novel *The Childhood of Bagrov, the Grandson* [6, 7]. Since Markov's time the study of complex chains has developed into one of the most powerful mathematical tools, applied far beyond quantitative linguistics in physics [8], chemistry [9], biology [10], finance [11], and even in numerical methods of estimation and optimization [12] and Google's PageRank algorithm [13].

To study correlation in structured processes, we take an information-theoretic view of Markovian complexity arising from temporal interdependency between observed symbols that are functions of chain states. That is, individual observation symbols are not themselves the chain states and therefore need not encode all the past. Specifically, we consider stationary, ergodic processes generated by hidden Markov chains (HMCs); introduced in the mid-twentieth century, as a generalization of Markov's complex chains, to model processes generated by communication channels [14]. When are these hidden processes described by finite Markov chains? When are they not Markovian? What's the informational signature in this case? And, what are "states" in the first place? Can we discover them from observations of a hidden process?

The following is the first in a series that addresses these questions: which have been answered, which can be answered, and which are open. Here, we concentrate on how the present—a sequence of $\ell$ consecutive measurements—statistically shields the past from the future, as a function of $\ell$. We introduce the *elusivity* $\sigma_\mu^\ell$ as a quantitative measure of the present failing to encode the past—a quantitative signature of Markov failure. We show how to calculate it explicitly via a novel construction and then describe and interpret its behavior through examples. The methods introduced also lead to compact expressions

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and efficient estimation of related measures—ephemeral, bound, and enigmatic informations—whose behaviors we also explore. As an application we use the results to reinterpret the persistent mutual information introduced by Ref. [15] as a measure of "emergence" in complex systems. Finally, we note that the sequel [16] is analytical, giving closed-form solutions and proving various properties, including several of those used here.

To address Markov's notion of complex chains, the next section reviews the minimal necessary background: measures of information content and correlation from information theory [17], their application to stochastic processes via computational mechanics [18, 19], and a recent analysis of the information content of the single time-step present within the context of the past and future [20]. This then sets the stage for a thorough analysis of Markovian complexity: Generalizing the previous framework so that the present can be an arbitrary duration. This gives rise to our main new result: expressing the elusivity in terms of a process's causal states. Notably, this result draws on the prior introduction of stochastic process models—the so-called bi-machines—that are agnostic with respect to the direction of time. We show how this leads to a simple and efficient expression for the elusive information and its companion measures. Those expressions, in turn, give the basis for further analytical development and for empirical estimation. With the general theory laid out, we make the ideas and methods concrete by calculating the elusivity and related quantities for a number of prototype complex processes, characterizing the variety of their convergence behaviors. Finally, we apply the insights gained to evaluate a proposed information-theoretic measure of emergence. We close by recapping the results and drawing conclusions for future applications.

## II. INFORMATION IN COMPLEX PROCESSES

### A. Processes

We are interested in a general stochastic *process* $\mathcal{P}$: the distribution of all of a system's behaviors or realizations $\{\ldots x_{-2}, x_{-1}, x_0, x_1, \ldots\}$ as specified by their joint probabilities $\Pr(\ldots X_{-2}, X_{-1}, X_0, X_1, \ldots)$. $X_t$ is a random variable that is the outcome of the measurement at the time $t$, taking values $x_t$ from a finite set $\mathcal{A}$ of all possible events. We denote a contiguous chain of random variables as $X_{0:\ell} = X_0 X_1 \cdots X_{\ell-1}$. Left indices are inclusive; right, exclusive. We suppress indices that are infinite. We consider only *stationary* processes for which $\Pr(X_{t:t+\ell}) = \Pr(X_{0:\ell})$ for all $t$ and $\ell$.

Our particular emphasis in the following is that a process $\Pr(X_{:0}, X_{0:\ell}, X_{\ell:})$ is a communication channel that transfers information from the *past* $X_{:0} = \ldots X_{-3}X_{-2}X_{-1}$ to the *future* $X_{\ell:} = X_\ell X_{\ell+1} X_{\ell+2} \ldots$ by storing parts of it in the *present* $X_{0:\ell} = X_0 X_1 \ldots X_{\ell-1}$ of length $\ell$. Of primary concern is whether $X_{:0} \to X_{0:\ell} \to X_{\ell:}$ forms a Markov chain in the sense of Ref. [21]:

$$\Pr(X_{-m:0}, X_{0:\ell}, X_{\ell:n}) = \Pr(X_{-m:0}|X_{0:\ell}) \Pr(X_{\ell:n}|X_{0:\ell}) \Pr(X_{0:\ell}) \ ,$$

for all $m, n \in \mathbb{Z}^+$.

### B. Channel Information

In analyzing this channel we need to measure a variety of information metrics, each capturing a different aspect of the information being communicated. The simplest is *Shannon entropy* [21]:

$$\mathrm{H}[X] = -\sum_{x \in \mathcal{X}} \Pr(x) \log_2 \Pr(x) \ . \tag{1}$$

Three other information-theoretic measures based on the entropy will be employed throughout. First, the *conditional entropy*, measuring the amount of information remaining in a variable $X$ (alphabet $\mathcal{X}$) once the information in a variable $Y$ (alphabet $\mathcal{Y}$) is accounted for:

$$\mathrm{H}[X|Y] = -\sum_{\substack{x \in \mathcal{X} \\ y \in \mathcal{Y}}} \Pr(x, y) \log_2 \Pr(x|y) \ . \tag{2}$$

Second, the deficiency of the conditional entropy relative to the full entropy is known as the *mutual information*, characterizing the information that is contained in both $X$ and $Y$:

$$\begin{aligned} \mathrm{I}[X:Y] &= \mathrm{H}[X] - \mathrm{H}[X|Y] \\ &= \sum_{\substack{x \in \mathcal{X} \\ y \in \mathcal{Y}}} \Pr(x, y) \log_2 \frac{\Pr(x, y)}{\Pr(x) \Pr(y)} \ . \end{aligned} \tag{3}$$

Last, we have the *conditional mutual information*, the mutual information between two variables once the information in a third ($Z$ with alphabet $\mathcal{Z}$) has been accounted for:

$$\mathrm{I}[X:Y|Z] = \sum_{\substack{x \in \mathcal{X} \\ y \in \mathcal{Y} \\ z \in \mathcal{Z}}} \Pr(x, y, z) \log_2 \frac{\Pr(x, y|z)}{\Pr(x|z) \Pr(y|z)} \ . \tag{4}$$

Perhaps the most naïve way of information-theoretically analyzing a process, capturing the randomness and dependencies in sequences of random variables, is via the

*block entropies*:

$$\mathrm{H}(\ell) = \mathrm{H}[X_{0:\ell}] \ . \tag{5}$$

This quantifies the amount of information in a contiguous block of observations. Its growth with $\ell$ gives insight into a process's randomness and structure [18, 22]:

$$\mathrm{H}(\ell) \approx \mathbf{E} + h_\mu \ell, \qquad \ell \gg 1 \ . \tag{6}$$

The asymptotic growth $h_\mu$, here, is a process's rate of information generation, or the *Shannon entropy rate*:

$$h_\mu = \mathrm{H}[X_0|X_{:0}] \ , \tag{7}$$

where the subscript $\mu$ denotes the specific probability measure over bi-infinite strings, defining the process of interest and its sequence probabilities. Lastly, the amount of future information predictable from the past is the past-future mutual information or *excess entropy*:

$$\begin{aligned} \mathbf{E} &= \mathrm{I}[X_{:0} : X_{0:}] \\ &= \mathrm{H}[X_{:0}, X_{0:}] - \mathrm{H}[X_{:0}|X_{0:}] \ . \end{aligned} \tag{8}$$

The excess entropy naturally arises when considering channels with a length $\ell = 0$ present, where it is effectively the only direct information quantity over the variables $X_{:0}$ and $X_{0:}$. It is well known that if the excess entropy vanishes, then there is no information temporally communicated by the channel [18].

Generically, Eq. (8) is of the form $\infty - \infty$, which is meaningless. In such situations one refers to finite sequences and then takes a limit:

$$\lim_{m,n\to\infty} \left( \mathrm{H}[X_{-m:0}, X_{0:n}] - \mathrm{H}[X_{-m:0}|X_{0:n}] \right) \ .$$

Here, we generally use the informal infinite variables in equations for clarity and simplicity unless the details of the limit are important for the analysis at hand. To be concrete, we write $f(X_{:0})$ to mean $\lim_{m\to\infty} f(X_{-m:0})$ and $f(X_{\ell:})$ to mean $\lim_{n\to\infty} f(X_{\ell:n})$.

### C. Information Atoms

The foregoing set up views a process as a channel that communicates the past to the future via the present. Our goal, then, is to analyze the channel's properties as a function of the present's length $\ell$. The cases of $\ell = 0$ and $\ell = 1$ have already been addressed: $\ell = 0$ in Ref. [23] and $\ell = 1$ in Ref. [20]. Our initial development closely mirror theirs. We borrow notation, but must include a superscript to denote the $\ell$-dependence of the quantities.

Broadly, our approach is to decompose the information in a random variable—such as the present—into components (atoms) associated with other random variables[1]. Our immediate concern is that of monitoring the amount of dependency remaining between the past and future if the present is known. We use the mutual information between the past and the future conditioned on the present to do so—the elusivity that will soon become our focus:

$$\begin{aligned} \sigma_\mu^\ell &= \mathrm{I}[X_{:0} : X_{\ell:}|X_{0:\ell}] \\ &= \mathrm{H}[X_{:0}|X_{0:\ell}] + \mathrm{H}[X_{\ell:}|X_{0:\ell}] - \mathrm{H}[X_{:0}, X_{\ell:}|X_{0:\ell}] \ . \end{aligned} \tag{9}$$

Note that $\sigma_\mu^0 = \mathbf{E}$.

Next, extending Ref. [20], we decompose the length-$\ell$ present. When considering only the past, the information in the present separates into two components: $\rho_\mu^\ell = \mathrm{I}[X_{:0} : X_{0:\ell}]$, the information that can be anticipated from the past, and $h_\mu^\ell = \mathrm{H}[X_{0:\ell}|X_{:0}]$, the random component that cannot be anticipated. Naturally, $\mathrm{H}[X_{0:\ell}] = h_\mu^\ell + \rho_\mu^\ell$. Connecting directly to Ref. [20], our $\rho_\mu^1$ is their $\rho_\mu$ and, likewise, our $h_\mu^1$ is their $h_\mu$.

If one also accounts for the future's behavior, then the random, unanticipated component $h_\mu^\ell$ breaks into two kinds of information: one part $b_\mu^\ell = \mathrm{I}[X_{0:\ell} : X_{\ell:}|X_{:0}]$ that, while a degree of randomness, is relevant for predicting the future; and the remaining part $r_\mu^\ell = \mathrm{H}[X_{0:\ell}|X_{:0}, X_{\ell:}]$ is ephemeral, existing only fleetingly in the present and then dissipating, leaving no trace on future behavior.

The redundant portion $\rho_\mu^\ell$ of $\mathrm{H}[X_{0:\ell}]$ itself splits into two pieces. The first part, $\mathrm{I}[X_{:0} : X_{0:\ell}|X_{\ell:}]$—also $b_\mu^\ell$ when the process is stationary—is shared between the past and the current observation, but its relevance stops there. The second piece $q_\mu^\ell = \mathrm{I}[X_{:0} : X_{0:\ell} : X_{\ell:}]$ is anticipated by the past, is present currently, and also plays a role in future behavior. Notably, this informational piece can be negative [20, 25].

Due to a duality between set-theoretic and information-theoretic operators, we can graphically represent the relationship between these various informations in a Venn-like display called an *information diagram* [26]; see Fig. 1. Similar to a Venn diagram, size indicates Shannon entropy rather than set cardinality and overlaps are not set intersection, but mutual information. Each area on the diagram represents one or another of Shannon's information measures.

As mentioned above, the past splits $\mathrm{H}[X_{0:\ell}]$ yielding two pieces: $h_\mu^\ell$, the part outside the past, and $\rho_\mu^\ell$, the part

[1] It is important to note that we are only associating portions of a random variable's information with other random variables. It is generically not possible to actually partition a random variable in to several other random variables, each of which has an entropy equal to the atom of interest [24].

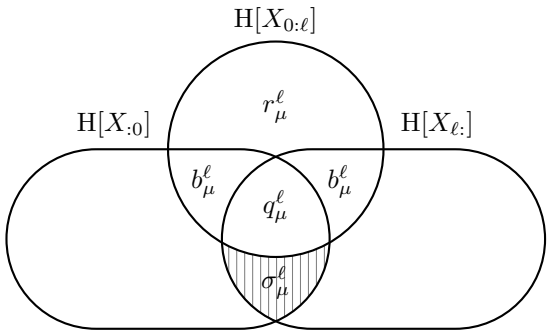
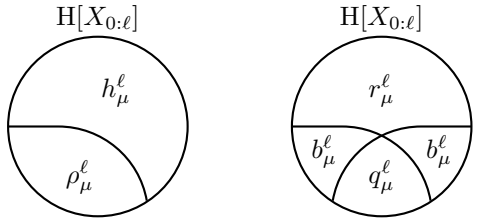



FIG. 1. The process information diagram that places the present in its temporal context: the past ($X_{:0}$) and the future ($X_{\ell:}$) partition the present ($X_{0:\ell}$) into four components with quantities $r_\mu^\ell$, $q_\mu^\ell$, and two with $b_\mu^\ell$. Notably, the component $\sigma_\mu^\ell$, quantifying the hidden dependency shared by the past and the future, lies outside of the present and so is *not* part of it.




(a) Decomposition due to the past.

(b) Decomposition due to the past and the future.

FIG. 2. Alternative decompositions of the present information H[$X_{0:\ell}$].

inside. This partitioning arises naturally when predicting a process [20]. To emphasize, Fig. 2a displays this decomposition. If we include the future in the diagram, we obtain a more detailed understanding of how information is transmitted from the past to the future. The past and the future together divide the present H[$X_{0:\ell}$] into four parts, as shown in Fig. 2b.

The process information diagram makes it rather transparent in which sense $r_\mu^\ell$ is an amount of *ephemeral information*: its information lies outside both the past and future and so it exists only in the present moment. It has no repercussions for the future and is no consequence of the past. It is the amount of information in the present observation neither communicated to the future nor from the past. With $\ell = 1$, this has been referred as the residual entropy rate [27], as it is the amount of uncertainty that remains in the present even after accounting for every other variable in the time series. It has also been studied as the erasure information [28] (there $H^-$), as it is the information irrecoverably erased in a binary erasure channel.

The *bound information* $b_\mu^\ell$ is the amount of spontaneously generated information present now, not explained by the past, but that has consequences for the future. In this sense it hints at being a measure of structural complexity [20, 27], though we discuss more direct measures of structure shortly.

Due to stationarity, the mutual information $\mathrm{I}[X_{0:\ell} : X_{\ell:}|X_{:0}]$ between the present $X_{0:\ell}$ and the future $X_{\ell:}$ conditioned on the past $X_{:0}$ is the same as the mutual information $\mathrm{I}[X_{0:\ell} : X_{:0}|X_{\ell:}]$ between the present $X_{0:\ell}$ and the past $X_{:0}$ conditioned on the future $X_{\ell:}$. Therefore they are both of size $b_\mu^\ell$, as shown in Fig. 1. This lends a symmetry to the process information diagram that need not exist for nonstationary processes.

### D. Elusivity

Two components remain in the process information diagram—two that have not been significantly analyzed previously. The first is $q_\mu^\ell = \mathrm{I}[X_{:0} : X_{0:\ell} : X_{\ell:}]$—the three-way mutual information (or *co-information* [25]) shared by the past, present, and future. Notably, unlike Shannon entropies and two-way mutual information, $q_\mu^\ell$ (and co-informations in general) can be negative. The other component $\sigma_\mu^\ell = \mathrm{I}[X_{:0} : X_{\ell:}|X_{0:\ell}]$, the quantity of primary interest here (shaded in Fig. 1) is the information shared between the past and the future that does not exist in the present. Since it measures dependency hidden from the present, we call it the *elusive information* or *elusivity* for short. Generally, it indicates that a process has hidden structures that are not appropriately captured by the present—that is, by finite random-variable blocks. In this case, and as we discuss at length towards the end, one must build models whose elements, which we call "states" below, represent how a process's internal mechanism is organized.

A process's internal organization somehow must store all the information from the past that is relevant for generating the future behavior. Only when the observed process is Markovian is it sufficient to keep the track of just the current observable or block of observables. For the general case of non-Markovian processes, though, information relevant for prediction is spread arbitrarily far back in the process's history and so cannot be captured by the present regardless of its duration. This fact is reflected in the existence of $\sigma_\mu^\ell$. When $\sigma_\mu^\ell > 0$ for all $\ell$, the description of the process requires determining its internal organization. This is one reason to build a model of the mechanism that generates sequences rather than simply describe a process as a list of sequences.

There are two basic properties that indicate the elusivity's importance. The first is that $\sigma_\mu^\ell$ decreases monotonically as a function of the present's length $\ell$. That

is, dependency cannot increase if we interpolate more random variables between the past and future.

**Proposition 1.** *For $k > 0$, $\sigma_\mu^\ell \geq \sigma_\mu^{\ell+k}$.*

**Proof.**

$$\begin{aligned}\sigma_\mu^\ell &= \mathrm{I}[X_{:0} : X_{\ell:}|X_{0:\ell}] \\ &= \mathrm{I}[X_{:0} : (X_{\ell:\ell+k}, X_{\ell+k:})|X_{0:\ell}] \\ &= \mathrm{I}[X_{:0} : X_{\ell:\ell+k}|X_{0:\ell}] + \mathrm{I}[X_{:0} : X_{\ell+k:}|X_{0:\ell}, X_{\ell:\ell+k}] \\ &= \mathrm{I}[X_{:0} : X_{\ell:\ell+k}|X_{0:\ell}] + \sigma_\mu^{\ell+k} \ ,\end{aligned}$$

*and by the non-negativity of conditional mutual informations [17], $\sigma_\mu^\ell \geq \sigma_\mu^{\ell+k}$.* □

The second property is that $\sigma_\mu^\ell$ indicates how poorly the present $X_{0:\ell}$ shields the past $X_{:0}$ and future $X_{\ell:}$. When it does, they are conditionally independent, given the present, and $\sigma_\mu^\ell$ vanishes. Due to this, it can be used to detect a process's *Markov order* $R$: the smallest $R$ for which $\Pr(X_{0:}|X_{:0}) = \Pr(X_{0:}|X_{-R:0})$.

**Proposition 2.** $\sigma_\mu^\ell = 0 \iff \ell \geq R$.

**Proof.** *By the definition of the Markov order $R$, a length-$R$ block is a sufficient statistic [29] for $X_{:0}$ about $X_{0:}$:*

$$\Pr(X_{0:}|X_{-R:0}) = \Pr(X_{0:}|X_{:0}, X_{-R:0}) \ ,$$

*and therefore also obeys [17]:*

$$\mathrm{I}[X_{:0} : X_{0:}] = \mathrm{I}[X_{-R:0} : X_{0:}] \ .$$

*Hence:*

$$\begin{aligned}\mathrm{I}[X_{:0} : X_{0:}] &= \mathrm{I}[(X_{:-R}, X_{-R:0}) : X_{0:}] \\ &= \mathrm{I}[X_{-R:0} : X_{0:}] + \mathrm{I}[X_{:-R} : X_{0:}|X_{-R:0}]\end{aligned}$$

*implies:*

$$\mathrm{I}[X_{:-R} : X_{0:}|X_{-R:0}] = \sigma_\mu^R = 0 \ ,$$

*due to stationarity.* □

To calculate $\sigma_\mu^\ell$, recall its definition as a conditional mutual information:

$$\begin{aligned}\sigma_\mu^\ell &= \mathrm{I}[X_{:0} : X_{\ell:}|X_{0:\ell}] \\ &= \sum_{\substack{x_{:0}\in X_{:0} \\ x_{0:\ell}\in X_{0:\ell} \\ x_{\ell:}\in X_{\ell:}}} \Pr(x_{:}) \log_2 \frac{\Pr(x_{:0}, x_{\ell:}|x_{0:\ell})}{\Pr(x_{:0}|x_{0:\ell}) \Pr(x_{\ell:}|x_{0:\ell})} \ ,\end{aligned}$$

where we used the notational shorthand for the bi-infinite joint distribution $\Pr(x_{:}) = \Pr(x_{:0}, x_{0:\ell}, x_{\ell:})$.

Note that for an order-$R$ Markov process, if $\ell \geq R$ the past and the future are independent over range $R$ [22] and so $\Pr(X_{:0}, X_{\ell:}|X_{0:\ell}) = \Pr(X_{:0}|X_{0:\ell}) \Pr(X_{\ell:}|X_{:0})$. With this, it is clear that $\sigma_\mu^\ell$ vanishes in such cases. This property has been discussed in prior literature as well [30].

Anticipating the needs of our calculations later, we replace conditional distributions with the joint ones: $\Pr(x_{:0}, x_{\ell:}|x_{0:\ell}) = \Pr(x_{:})/\Pr(x_{0:\ell})$ and $\Pr(x_{:0}|x_{0:\ell}) = \Pr(x_{:0}, x_{0:\ell})/\Pr(x_{0:\ell})$, obtaining:

$$\sigma_\mu^\ell = \sum_{\substack{x_{:0}\in X_{:0} \\ x_{0:\ell}\in X_{0:\ell} \\ x_{\ell:}\in X_{\ell:}}} \Pr(x_{:}) \log_2 \frac{\Pr(x_{0:\ell}) \Pr(x_{:})}{\Pr(x_{:0}, x_{0:\ell}) \Pr(x_{0:\ell}, x_{\ell:})} \ . \quad (10)$$

Notably, all the terms needed to compute $\sigma_\mu^\ell$ are either $\Pr(x_{:0}, x_{0:\ell}, x_{\ell:})$ or marginals thereof. Our next goal, therefore, is to develop the theoretical infrastructure necessary to compute that distribution in closed form.

Similar expressions, which we use later on but do not record here, can be developed for the other information measures $h_\mu^\ell$, $r_\mu^\ell$, $b_\mu^\ell$, and $q_\mu^\ell$.

## III. STRUCTURAL COMPLEXITY

To analytically calculate the elusive information $\sigma_\mu^\ell$ and related measures we must go beyond the information theory of sequences and introduce *computational mechanics*, the theory of process structure [19]. The representation it uses for a given process is a form of *hidden Markov model* (HMM) [31]: the *$\epsilon$-machine*, which consists of a set $\boldsymbol{\mathcal{S}}$ of *causal states* and a transition dynamic $T$. $\epsilon$-Machines satisfy three conditions: irreducibly, unifilarity, and probabilistically distinct states [32]. *Irreducibly* implies that the associated state-transition graph is strongly connected. *Unifilarity*, perhaps the most distinguishing feature, means for each state $\sigma \in \boldsymbol{\mathcal{S}}$ and each observed symbol $x$ there is at most one outgoing transition from $\mathcal{S}$ labeled $x \in \mathcal{A}$. Critically, unifilarity enables one to directly calculate various process quantities, such as conditional mutual informations, using properties of the hidden (causal) states. Notably, many of these quantities cannot be directly calculated using the states of general (nonunifilar) HMMs. Finally, an HMM has *probabilistically distinct states* when, for every pair of states $\mathcal{S}$ and $\mathcal{S}'$, there exists a word $w$ such that the probability of observing $w$ from each state is distinct: $\Pr(w|\mathcal{S}) \neq \Pr(w|\mathcal{S}')$. An irreducible, unifilar model with probabilistically distinct states is minimal in the sense that no model with fewer states or transitions generates the process. An HMM satisfying these three properties is an $\epsilon$-machine.

### A. Constructing the $\epsilon$-Machine

Given a process, how does one construct it's $\epsilon$-machine? First, the set of a process's *forward causal states*:

$$\boldsymbol{\mathcal{S}}^+ = X_{:}/\sim_\epsilon^+ \tag{11}$$

is the partition defined via the *causal equivalence relation*:

$$x_{:t} \sim_\epsilon^+ x'_{:t} \Leftrightarrow \Pr(X_{t:}|X_{:t} = x_{:t}) = \Pr(X_{t:}|X_{:t} = x'_{:t}) \ . \tag{12}$$

That is, each causal state $\sigma^+ \in \boldsymbol{\mathcal{S}}^+$ is an element of the coarsest partition of a process's pasts such that every $x_{:0} \in \sigma^+$ makes the same prediction $\Pr(X_{0:}|\cdot)$. In fact, the causal states are the *minimal sufficient statistic* of the past to predict the future. We define the *reverse causal states*:

$$\boldsymbol{\mathcal{S}}_t^- = X_{t:}/\sim_\epsilon^- \ . \tag{13}$$

by similarly partitioning the process's futures:

$$x_{t:} \sim_\epsilon^- x'_{t:} \Leftrightarrow \Pr(X_{:t}|X_{t:} = x_{t:}) = \Pr(X_{:t}|X_{t:} = x'_{t:}) \ . \tag{14}$$

Second, the causal equivalence relation provides a natural unifilar dynamic over the states. For each state $\sigma$ and next symbol $x$, either there is a successor state $\sigma'$ such that the updated past $x_{:t+1} = x_{:t}x \in \sigma'$, for all $x_{:t} \in \sigma$, or $x_{:t+1}$ does not occur. Due to causal-state equivalence, every past within a state collectively either can or cannot be followed by a given symbol. Moreover, since the causal states form a partition of all pasts, there is at most one causal state to which each past can advance.

For a general HMM with states $\rho \in \mathcal{R}$, its *symbol-labeled transition matrix* $T^{(x)}$ has elements that give the probability of going from state $\rho$ to state $\rho'$ and generating the symbol $x$:

$$T^{(x)}_{\rho\rho'} \equiv \Pr(X_t = x, \mathcal{R}_{t+1} = \rho'|\mathcal{R}_t = \rho) \ . \tag{15}$$

Furthermore, the internal-state dynamics is governed by the stochastic matrix $T = \sum_x T^{(x)}$. Its unique left eigenvector $\pi$, associated with eigenvalue 1, gives the asymptotic state probability $\Pr(\rho)$. By extension, the transition matrix giving the probability of a word $w = x_0x_1 \cdots x_{\ell-1}$ of length $\ell$ is the product of transition matrices of each symbol in $w$:

$$\begin{aligned} T^{(w)} &\equiv \prod_{x_i \in w} T^{(x_i)} \\ &= T^{(x_0)}T^{(x_1)} \cdots T^{(x_{\ell-1})} \ . \end{aligned} \tag{16}$$

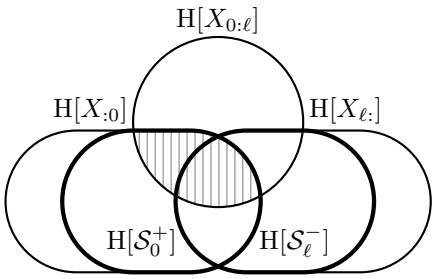



FIG. 3. Mutual information $\mathrm{I}[X_{:0} : X_{0:\ell}]$ between the past and the present (shaded) is equivalent to the mutual information $\mathrm{I}[\mathcal{S}_0^+ : X_{0:\ell}]$ between the forward causal state and the present.

### B. Rendering $\sigma_\mu^\ell$ Finitely Computable

We can put the forward and reverse causal states to use since they are proxies for a process's semi-infinite pasts and futures, respectively. See, e.g., Fig. 3. In this way, we transform Eq. (9) into a form containing only finite sets of random variables. We calculate directly:

$$\begin{aligned} \sigma_\mu^\ell &= \mathrm{I}[X_{:0} : X_{\ell:}|X_{0:\ell}] \\ &= \mathrm{I}[X_{:0} : (X_{0:\ell}, X_{\ell:})] - \mathrm{I}[X_{:0} : X_{0:\ell}] \\ &= \mathrm{I}[X_{:0} : X_{0:}] - \mathrm{I}[X_{:0} : X_{0:\ell}] \\ &\overset{(a)}{=} \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^-] - \mathrm{I}[\mathcal{S}_0^+ : X_{0:\ell}] \\ &= \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^-] - (\mathrm{I}[\mathcal{S}_0^+ : X_{0:\ell} : \mathcal{S}_0^-] + \mathrm{I}[\mathcal{S}_0^+ : X_{0:\ell}|\mathcal{S}_0^-]) \\ &\overset{(b)}{=} \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^-] - \mathrm{I}[\mathcal{S}_0^+ : X_{0:\ell} : \mathcal{S}_0^-] \\ &= \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^-|X_{0:\ell}] \\ &= \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^- : \mathcal{S}_\ell^-|X_{0:\ell}] + \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^-|X_{0:\ell}, \mathcal{S}_\ell^-] \\ &\overset{(c)}{=} \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_0^- : \mathcal{S}_\ell^-|X_{0:\ell}] \\ &= \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_\ell^-|X_{0:\ell}] - \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_\ell^-|X_{0:\ell}, \mathcal{S}_0^-] \\ &\overset{(d)}{=} \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_\ell^-|X_{0:\ell}] \\ &\quad - (\mathrm{H}[\mathcal{S}_0^+|X_{0:\ell}, \mathcal{S}_0^-] - \mathrm{H}[\mathcal{S}_0^+|\mathcal{S}_0^-, X_{0:\ell}, \mathcal{S}_\ell^-]) \\ &= \mathrm{I}[\mathcal{S}_0^+ : \mathcal{S}_\ell^-|X_{0:\ell}] \ . \end{aligned} \tag{17}$$

Above, $(a)$ is true due to Eqs. (12) to (13) and Ref. [33], $(b)$ is true due to Eqs. (13) and (14), $(c)$ is true due to Eq. (14) and unifilarity, and finally $(d)$ is true due to both entropy terms being equal to $\mathrm{H}[\mathcal{S}_0^+|\mathcal{S}_0^-]$ by Eqs. (13) and (14). That is, $\mathcal{S}_0^-$ informationally subsumes both $X_{0:\ell}$ and $\mathcal{S}_\ell^-$ when it comes to $X_{:0}$ and, therefore, also when it comes to $\mathcal{S}_0^+$. All other equalities are basic information identities found in Ref. [21].

In this way, Eq. (17) says that Eq. (10) becomes, in

terms of causal states, a new expression for elusivity:

$$\sigma_\mu^\ell = \sum_{\substack{\sigma_0^+ \in \mathcal{S}_0^+ \\ x_{0:\ell} \in X_{0:\ell} \\ \sigma_\ell^- \in \mathcal{S}_\ell^-}} \Pr(\sigma_0^+, x_{0:\ell}, \sigma_\ell^-) \log_2 \frac{\Pr(x_{0:\ell}) \Pr(\sigma_0^+, x_{0:\ell}, \sigma_\ell^-)}{\Pr(\sigma_0^+, x_{0:\ell}) \Pr(x_{0:\ell}, \sigma_\ell^-)} \ . \tag{18}$$

We transformed the key distribution $\Pr(x_{:0}, x_{0:\ell}, x_{\ell:})$ over random variables $X_{:0}$ and $X_{\ell:}$ with cardinality of the continuum to $\Pr(\sigma_0^+, x_{0:\ell}, \sigma_\ell^-)$ over $\mathcal{S}^+$ and $\mathcal{S}^-$ with typically smaller cardinality. When the causal states are finite or countably infinite, the benefit is substantial. We will now turn our attention to computing this joint distribution.

Since the distribution is over both forward and reverse causal states, we must track both simultaneously. The key tool for this is Ref. [34]'s bidirectional machine or *bimachine*. We point the reader there for details regarding their construction and properties. One feature we need immediately, though, is that bimachine states $\rho_t = (\sigma_t^+, \sigma_t^-)$ are pairs of forward and reverse causal states.

Generally, given an HMM with states $\rho \in \mathcal{R}$, we can construct the distribution of interest if we can find a way to build distributions of the form $\Pr(\rho_i, w, \rho_j)$: the probability of being in state $\rho_i$, generating the word $w$, and ending in state $\rho_j$. The word transition matrix (Eq. (16)) gives exactly this and allows us to build the distribution directly:

$$\Pr(\rho_i, w, \rho_j) = (\pi \circ \mathbf{1}_i) T^{(w)} \mathbf{1}_j^\intercal \ , \tag{19}$$

where $\rho_i$ and $\rho_j$ are the states of an arbitrary HMM, $a \circ b$ is the Hadamard (elementwise) product of vectors $a$ and $b$, and $\mathbf{1}_i$ is the row vector with all its elements zero except for the $i^{\text{th}}$, which is 1.

Applying Eq. (19) to the bimachine, we arrive at the distribution $\Pr((\mathcal{S}_0^+, \mathcal{S}_0^-), X_{0:\ell}, (\mathcal{S}_\ell^+, \mathcal{S}_\ell^-))$, which can be marginalized to $\Pr(\mathcal{S}_0^+, X_{0:\ell}, \mathcal{S}_\ell^-)$, the distribution needed to compute Eq. (18). Figure 4 illustrates this distribution for a present of length $\ell = 3$ in the setting of the process's random variable lattice and the forward and reverse causal state processes.

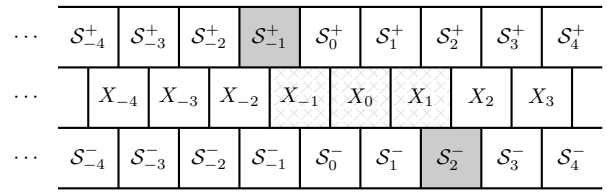



FIG. 4. Random variable lattice illustrating the relationship between forward causal states $\mathcal{S}_t^+$, observed symbols $X_t$, and reverse causal states $\mathcal{S}_t^-$. The variables in the distribution $\Pr(\mathcal{S}_{-1}^+, X_{-1:2}, \mathcal{S}_2^-)$ are highlighted. In particular, elusivity $\sigma_\mu^3$ is the mutual information between the two shaded cells ($\mathcal{S}_{-1}^+$ and $\mathcal{S}_2^-$) conditioned on the hatched cells ($X_{-1:2} = X_{-1}X_0X_1$).

### C. Companion Atomic Measures

Causal-state expressions for $h_\mu^\ell$, $r_\mu^\ell$, $b_\mu^\ell$, and $q_\mu^\ell$ that we use in the following are:

$$\begin{aligned} h_\mu^\ell &= \mathrm{H}[X_{0:\ell} | \mathcal{S}_0^+] \ , \\ r_\mu^\ell &= \mathrm{H}[X_{0:\ell} | \mathcal{S}_0^+ : \mathcal{S}_\ell^-] \ , \\ b_\mu^\ell &= \mathrm{I}[X_{0:\ell} : \mathcal{S}_0^+ | \mathcal{S}_\ell^-] \ , \text{ and} \\ q_\mu^\ell &= \mathrm{I}[\mathcal{S}_0^+ : X_{0:\ell} : \mathcal{S}_\ell^-] \ . \end{aligned}$$

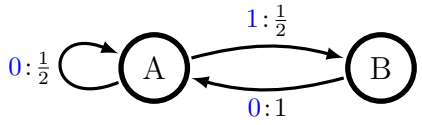
These are derived in ways paralleling that above for $\sigma_\mu^\ell$ and so we do not give detail. They, too, also depend on the joint distribution above in Eq. (19) and its marginals.

## IV. EXAMPLES

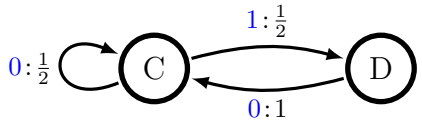
Let's consider several example processes, to illustrate calculation methods and to examine the behavior of $\sigma_\mu^\ell$ and companion measures.

### A. Golden Mean Process

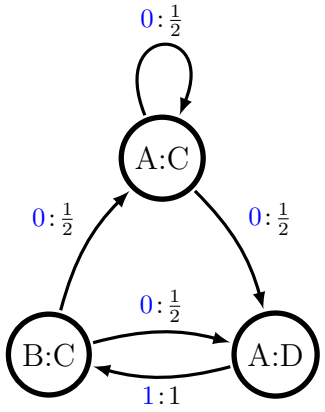
As the first example we analyze the Golden Mean (GM) Process, whose $\epsilon$-machines and bimachine state-transition diagrams are given in Fig. 5. The GM Process consists of all bi-infinite strings such that no consecutive 1s occur, with probabilities such that either symbol is equally likely following a 0. A stochastic generalization of subshifts of finite type [35] this process can be described by a Markov chain with order $R = 1$. Due to Prop. 2 we expect $\sigma_\mu^1 = 0$. To verify this, we compute each term of Eq. (18) using the edges of the bimachine, Fig. 5c, and the invariant




(a) Forward $\epsilon$-machine.




(b) Reverse $\epsilon$-machine.

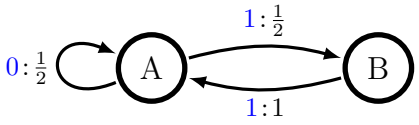



(c) Bimachine.

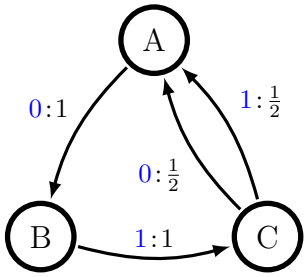
FIG. 5. The several faces of the Golden Mean (GM) Process. Each representation consists of states (labeled circles) and transitions (arrows) labeled "*symbol* : *probability*". The bimachine is effectively the Cartesian product of the forward and reverse machines, though constructed here for forward-time generation, and therefore is generically nonunifilar. For example, the state B:C means that the machine is in the superposition of forward state B and reverse state C. Going forward (rightward in Fig. 4) we know that state B must output a 0 and transition to state A. Being in reverse state C we must have transitioned to there on a 0 coming from either state C or state D. Thus, we have transitions from B:C to both A:C and A:D on a 0.

state distribution $\pi = \left(1/3, 1/3, 1/3\right)$:

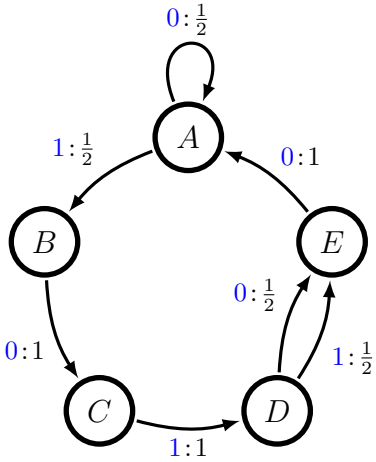
$$\frac{\Pr(0)\Pr(A,0,C)}{\Pr(A,0)\Pr(0,C)} = \frac{2/3 \cdot 1/6}{1/3 \cdot 1/3} = 1 \ ,$$
$$\frac{\Pr(0)\Pr(A,0,D)}{\Pr(A,0)\Pr(0,D)} = \frac{2/3 \cdot 1/6}{1/3 \cdot 1/3} = 1 \ ,$$
$$\frac{\Pr(0)\Pr(B,0,C)}{\Pr(B,0)\Pr(0,C)} = \frac{2/3 \cdot 1/6}{1/3 \cdot 1/3} = 1 \ ,$$
$$\frac{\Pr(0)\Pr(B,0,D)}{\Pr(B,0)\Pr(0,D)} = \frac{2/3 \cdot 1/6}{1/3 \cdot 1/3} = 1 \ , \text{ and}$$
$$\frac{\Pr(1)\Pr(A,1,C)}{\Pr(A,1)\Pr(1,C)} = \frac{1/3 \cdot 1/3}{1/3 \cdot 1/3} = 1 \ .$$




(a) The Even Process.




(b) The Noisy Period Three (NP3) Process.




(c) The Noisy Random Phase-Slip (NRPS) Process.

FIG. 6. $\epsilon$-Machines for the Example Processes.

We see that the argument of each $\log_2$ in Eq. (18) is 1, confirming that $\sigma_\mu^1 = 0$.

### B. Information Measures versus Present Length

We now investigate the behavior of $\sigma_\mu^\ell$ and its companions $q_\mu^\ell$, $b_\mu^\ell$, and $r_\mu^\ell$ for several example processes: the aforementioned GM, the Even, the Noisy Period-Three (NP3), and the Noisy Random Phase-Slip (NRPS) Processes. The $\epsilon$-machines for the latter are shown in Fig. 6. Each exhibits different convergence behaviors with $\ell$ for the differing measures; see the graphs in Fig. 7. We now turn to characterizing each of them.

We first consider $\sigma_\mu^\ell$, seen in Fig. 7's upper left panel. While for each process $\sigma_\mu^\ell$ vanishes with increasing $\ell$, the convergence behaviors differ. The Golden Mean Process is identically zero at all lengths due to its order-1 Markov

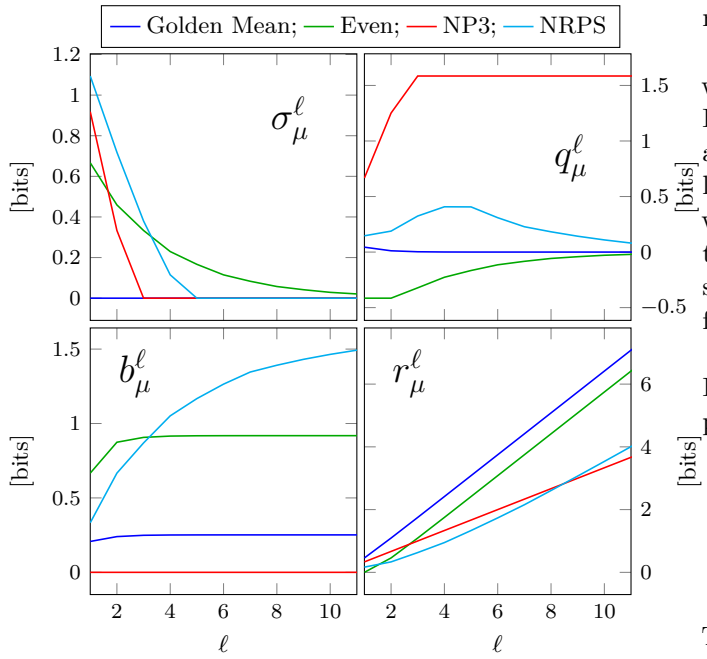



FIG. 7. Information measures as a function of the present's length $\ell$. Since the examples are stationary, finitary processes, both $\sigma_\mu^\ell$ and $q_\mu^\ell$ converge to zero with increasing $\ell$ and $b_\mu^\ell$ converges to a constant value of $\mathbf{E}$. And so, the growth of $\mathrm{H}(\ell)$ is entirely captured in $r_\mu^\ell$, and it grows linearly with $\ell$.

nature, just noted. The NRPS Process, with a Markov order of $R = 5$, has nonzero $\sigma_\mu^\ell$ for $\ell < 5$ and zero $\sigma_\mu^\ell = 0$ beyond. Finally, both the Even and Nemo Processes are infinite-order Markov and so their $\sigma_\mu^\ell$ never exactly vanishes, though they converge exponentially fast. The next section, Section IV C, analyzes exponential convergence in more detail.

Next, consider $q_\mu^\ell$ and $b_\mu^\ell$ which, as it turns out, are closely associated. To see why, first examine the large-$\ell$ limit:

$$\begin{aligned}\lim_{\ell\to\infty} \mathrm{I}[X_{0:\ell} : X_{\ell:}] &= \mathbf{E} \\ &= q_\mu^\infty + b_\mu^\infty \ .\end{aligned}$$

Second, we can decompose stationary, finite-$\mathbf{E}$ processes into two classes: those with state mixing and those without. *State mixing* refers to the convergence of initial state distributions to a unique invariant distribution, one that in particular does not oscillate. The GMP, Even, and NRPS Processes are examples of those with state mixing, while the NP3 Process asymptotic state distribution is period-3, when starting from typical initial distributions. With state mixing $X_{:0}$ and $X_{\ell:}$ become independent in the infinite-$\ell$ limit, and so $q_\mu^\infty = 0$ and we conclude $b_\mu^\infty = \mathbf{E}$. That is, the entire contribution to excess entropy comes from the bound information. Without state mixing, though, $b_\mu^\ell = 0$, and so $q_\mu^\infty = \mathbf{E}$.

These two classes of convergence behavior are apparent when comparing Fig. 7's upper-right and lower-left panels. In the upper-right, $q_\mu^\ell$ converges to zero for the GM, Even, and NRPS Processes. The NP3 Process, in contrast, limits to a constant value: it's excess entropy $\mathbf{E} = \log_2 p$, where $p = 3$ is the period of the internal state cycle. In the lower-left, $b_\mu^\ell$ limits to constant values for the three state-mixing processes, while the NP3 Process $b_\mu^\ell$ is zero for all $\ell$.

Finally, the ephemeral information $r_\mu^\ell$, plotted in Fig. 7's lower-right panel, also depends on whether a process is state mixing or not. If it is, then:

$$\begin{aligned}\mathrm{H}(\ell) &= \mathbf{E} + \ell h_\mu \\ &= q_\mu^\ell + 2b_\mu^\ell + r_\mu^\ell \\ &= 0 + 2\mathbf{E} + r_\mu^\ell \ .\end{aligned}$$

That is, $r_\mu^\ell = -\mathbf{E} + \ell h_\mu$. Though, in the case of no state mixing, we have:

$$\begin{aligned}\mathrm{H}(\ell) &= q_\mu^\ell + 2b_\mu^\ell + r_\mu^\ell \\ &= \mathbf{E} + 0 + r_\mu^\ell \ .\end{aligned}$$

That is, $r_\mu^\ell = \ell h_\mu$. So, in either case, $r_\mu^\ell$ grows linearly asymptotically with a rate of $h_\mu$. If there is state mixing, however, it has a subextensive part equal to $-\mathbf{E}$.

### C. Exponential Convergence of $\sigma_\mu^\ell$

One way to classify processes is whether or not an observer can determine the causal state a process is in from finite or infinite sequence measurements. If so, then the process is *synchronizable*. All of the previous examples are synchronizable. References [36, 37] proved that for any synchronizable process described by a finite-state HMM, there exist constants $K > 0$ and $0 < \alpha < 1$ such that:

$$h_\mu(\ell) - h_\mu \leq K\alpha^\ell, \quad \textit{for all } \ell \in \mathbb{N} \ , \tag{20}$$

where:

$$h_\mu(\ell) = \mathrm{H}(\ell) - \mathrm{H}(\ell - 1) \ .$$

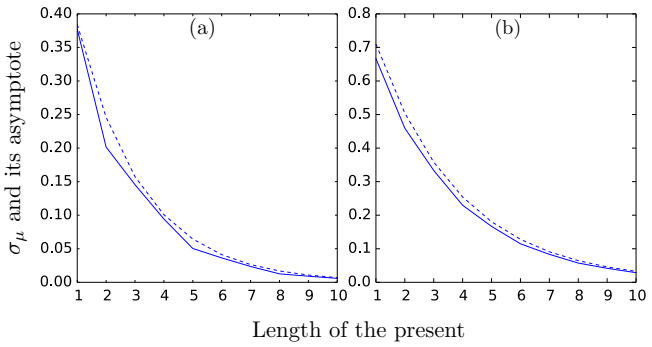



FIG. 8. $\sigma_\mu^\ell$ (solid) and its asymptote (dashed) for (a) Nemo Process with $K = 0.6$ and $\alpha = 0.64$ and (b) Even Process with $K = 1$ and $\alpha = 0.71$.

Note that $h_\mu(1) = \rho_\mu^1$. One well known identity [18] is that the sum of the $h_\mu(\ell)$ terms is the excess entropy:

$$\mathbf{E} = \sum_{k=1}^{\infty} (h_\mu(k) - h_\mu) \tag{21}$$

$$= \rho_\mu^1 + \sum_{k=2}^{\infty} (h_\mu(k) - h_\mu) \ . \tag{22}$$

This provides a new identity [16]:

$$\sigma_\mu^1 = \sum_{k=2}^{\infty} (h_\mu(k) - h_\mu) \ , \tag{23}$$

which can be generalized to:

$$\sigma_\mu^\ell = \sum_{k=\ell+1}^{\infty} (h_\mu(k) - h_\mu) \ . \tag{24}$$

Applying the bound from Eq. (20) to each term, we find:

$$\sum_{k=\ell+1}^{\infty} \left(h_\mu(k) - h_\mu\right) \leq \sum_{k=\ell+1}^{\infty} K\alpha^k \ .$$

The right-hand side, being a convergent geometric series, yields:

$$\sigma_\mu^\ell \leq \frac{K\alpha^{\ell+1}}{1-\alpha} \ ,$$

or simply:

$$\sigma_\mu^\ell \leq K'\alpha^\ell \ . \tag{25}$$

We now drop the prime, simplifying the notation. Thus, the elusive information vanishes exponentially fast for synchronizable processes.

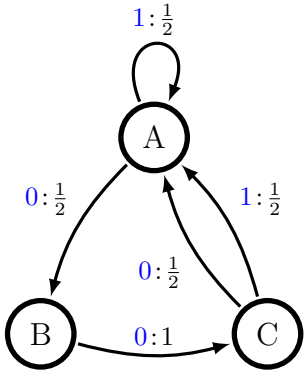



FIG. 9. The Nemo Process.

Figure 8 compares $\sigma_\mu^\ell$ with its best-fit exponential bound for two different processes: the Nemo Process shown in Fig. 9 and the Even Process. For each, the solid line is $\sigma_\mu^\ell$ and the dashed is the fit. Estimated values for the Nemo Process are $K = 0.6$ and $\alpha = 0.64$. The fit parameters for the Even Process are $K = 1.0$ and $\alpha = 0.71$. They were estimated in accordance with the conditions stated for Eq. (20). The fits validate the convergence in Eq. (25).

## V. MEASURES OF EMERGENCE?

The elusive information $\sigma_\mu^\ell$ conditions on a present of length $\ell$. What if we do not condition, simply ignoring the present? It becomes the *persistent mutual information* (PMI) [15, 38]:

$$\mathrm{PMI}(\ell) = \mathrm{I}[X_{:0} : X_{\ell:}] \ . \tag{26}$$

Notably, PMI($\infty$) was offered up as a measure of "emergence" in general complex systems. Our preceding analysis, though, gives a more nuanced view of this interpretation, especially when emergence is considered in light of structural criteria introduced previously [39, 40]. Our framework reveals that PMI($\ell$) is not an atomic measure; rather it consists of two now-familiar components:

$$\mathrm{PMI}(\ell) = q_\mu^\ell + \sigma_\mu^\ell \ . \tag{27}$$

Which component is most important? Are both? Which is associated with emergence? Both?

Section IV C showed that synchronizable processes have $\sigma_\mu^\ell \to 0$. So, for this broad class at least, PMI($\infty$) $= q_\mu^\infty$. Based on extensive process surveys that we do not report on here, we conjecture that $\sigma_\mu^\infty = 0$ holds even more generally. And so, it appears that PMI($\infty$) generally is dominated by the multivariate mutual information $q_\mu^\infty$.

Moreover, recalling the analysis of $q_\mu^\ell$ for the Noisy Period-3 Process shown in the upper-right panel of Fig. 7, it appears that PMI($\infty$) is only sensitive to periodicity or noisy periodicity, giving $\log_2 p$, where $p$ is the period.

As a test of our conjecture that the elusive information vanishes and that PMI($\infty$) is dominated by $q_\mu^\infty$, we applied our information-measure estimation methods to the symbolic dynamics generated by the Logistic Map of the unit interval as a function of its control parameter $r$. Figure 10 plots the results. Indeed, the elusive information does vanish. Thus, we conclude that PMI($\infty$) is a property of $q_\mu^\infty$.

In addition, our simulation results reproduced those in Ref. [15]'s PMI($\infty$) analysis of the Logistic Map; though, their estimation method for PMI($\infty$) differs markedly. Here, we calculate via the Logistic Map symbolic dynamics; there, joint distributions over the continuous unit-interval domain were used. Both investigations lead to the conclusion that PMI($\infty$) is equal to (the logarithm of) the number of chaotic "bands" cyclically permuted or the period of the periodic orbit at a given parameter value. In short, PMI($\infty$) is a measure of *non-mixing* dynamics.

Given the restricted form of structure (periodicity) to which it is sensitive, PMI($\infty$) cannot be taken as a general measure for detecting the emergence of organization in complex systems. No matter, though a quarter of a century old, the statistical complexity [41]—a direct measure of structural organization and stored information—continues to fill the role of detecting emergent organization quite well. Moreover, computational mechanics' $\epsilon$-machines directly show what the emergent organization is.

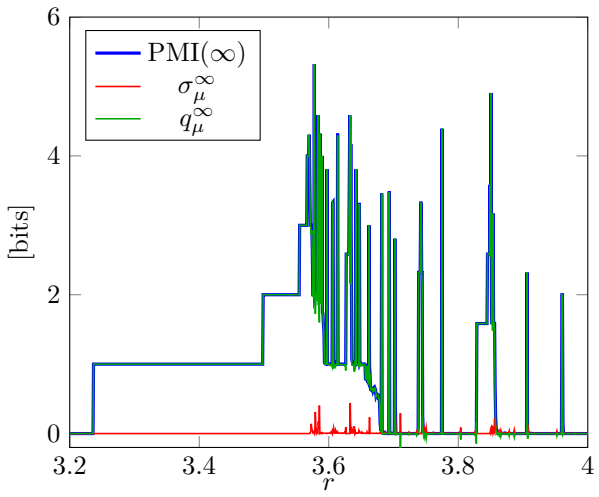



FIG. 10. Persistent mutual information PMI($\infty$), elusivity $\sigma_\mu^\infty$, and multivariate mutual information $q_\mu^\infty$ of the Logistic Map symbolic dynamics as a function of map control parameter $r$. Recall that the symbolic dynamics does not see period-doubling until $r$ is above the appearance of the associated superstable periodic orbit. This discrepancy in the appearance of periodicity as a function of $r$ does not occur when the map is chaotic. (Cf. Fig. 1 of Ref. [15].)

## VI. CONCLUSION

We first defined the elusive information and developed a closed-form analytic expression to calculate it from a process's hidden Markov model. The sequel [16] shows how to use spectral methods [42] to develop alternative closed-form expressions for the elusive information and its companions, giving exact expressions and a direct understanding of the origin of their convergence behaviors.

Investigating how the present shields the past and future is essentially a study of what Markov order means for structured processes. It gives much insight into the process of modeling building and even into general concerns about the emergence of organization in complex systems. First, this study of Markovian complexity gives a common ground on which to contrast structural inference and emergence, showing that we should not conflate these two distinct questions. Second and perhaps most constructively, though, it sheds light on the challenges of inference for complex systems. In particular, when $\sigma_\mu^\ell > 0$ sequence statistics are inadequate for modeling and so we must employ state-based models to properly, finitely represent a process's internal organization.

The results show rather directly how present observables typically do not contain all of the information that correlates the past and the future. One consequence is that instantaneous measurements are not enough. This means, exactly, that Markov chain models of complex physical systems are fundamentally inadequate, though eminently helpful and simplifying when they are appropriate representations. When predicting the behavior and structure of complex systems, the larger consequence is that we must build state-based models and not use mere look-up tables or sequence histograms. That said, states are little more than a conditional coarse-graining of sequences into subsets that are more compactly predictive than sequences alone. And this means, in turn, that monitoring *only* prediction performance is inadequate. We must also monitor model complexity, not as an antidote to over-fitting, but as a fundamental goal for both prediction and understanding hidden mechanisms. This, we believe, most fully respects Markov's contribution to the sciences.

## ACKNOWLEDGMENTS

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