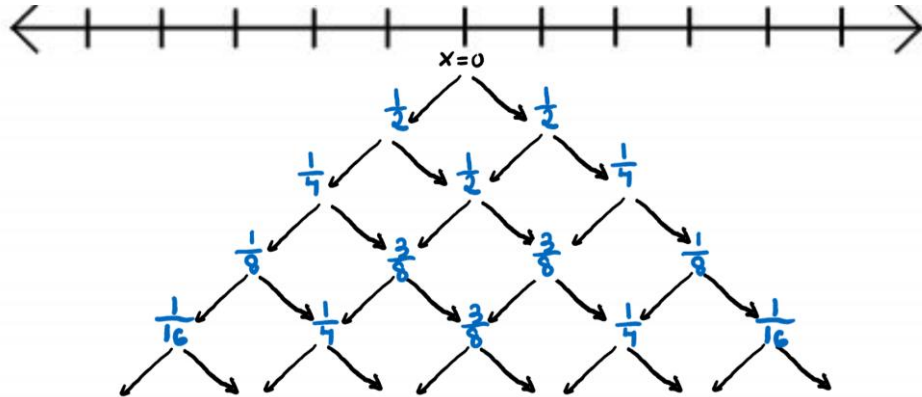


# Two Particle Quantum Walk

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What does drunkenness look like in the Quantum Realm?

# Random Walk

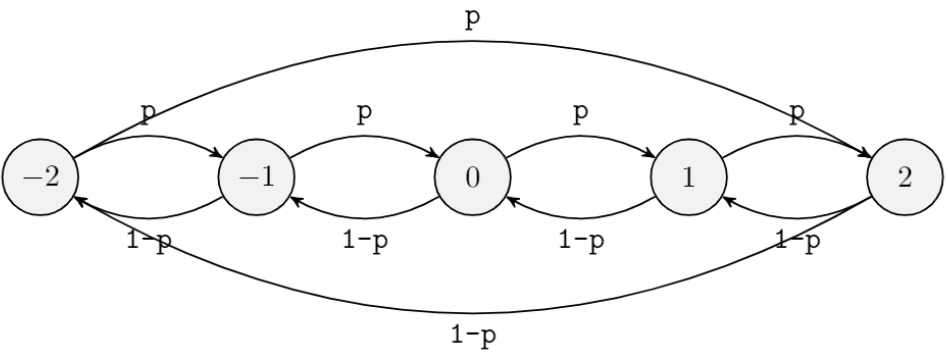


Random walk is a stochastic process, that describes a path that consists of a succession of random steps on some mathematical space.

A drunkard's walk—where the choice of whether to step to the right or the left is made randomly by the toss of a coin

# Random Walk as a Markov Chain

- Random walk on a  $N$  vertex graph is given by an  $N \times N$  matrix -  $T$  where,  $T_{ij}$  represents the probability of making a transition from  $i$  to  $j$  and  $\sum_{j=1}^N T_{ij} = 1$
- The state space for a  $1D$  Random Walk:  $\{0, \pm 1, \pm 2\}$  with periodic boundary conditions



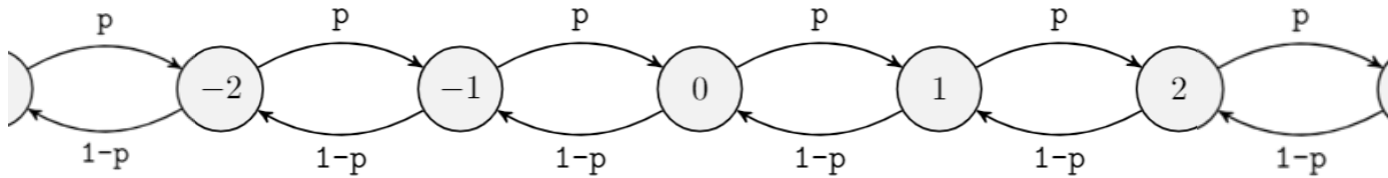
$$T = \begin{bmatrix} 0 & p & 0 & 0 & 1-p \\ 1-p & 0 & p & 0 & 0 \\ 0 & 1-p & 0 & p & 0 \\ 0 & 0 & 1-p & 0 & p \\ p & 0 & 0 & 1-p & 0 \end{bmatrix} = T^{(L)} + T^{(R)}$$

$$T^{(R)} = p \begin{bmatrix} 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 \end{bmatrix} \quad T^{(L)} = (1-p) \begin{bmatrix} 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \end{bmatrix}$$

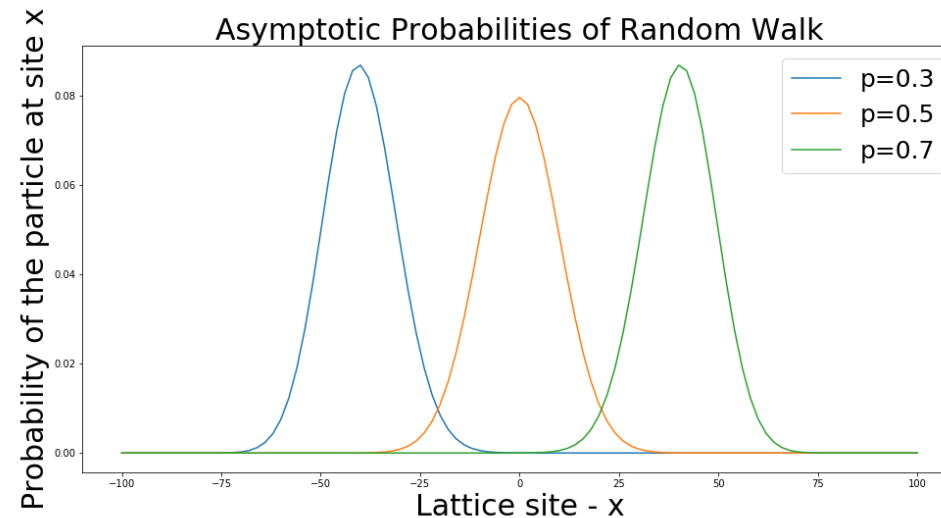
- $\mu(t + 1) = \mu(t)T$  , where  $\mu \in \mathbb{R}^N$
- $T$  can be used to obtain the asymptotic probability distribution  $\pi$
- Here  $\pi = \{\frac{1}{5}, \frac{1}{5}, \frac{1}{5}, \frac{1}{5}, \frac{1}{5}\}$  , when  $p = \frac{1}{2}$  , where  $\pi \in \mathbb{R}^N$

# Random Walk on infinite 1D chain

- For an infinite 1D chain :
- The state space for infinite 1D Random Walk:  $\{0, \pm 1, \pm 2, \pm 3 \dots\}$

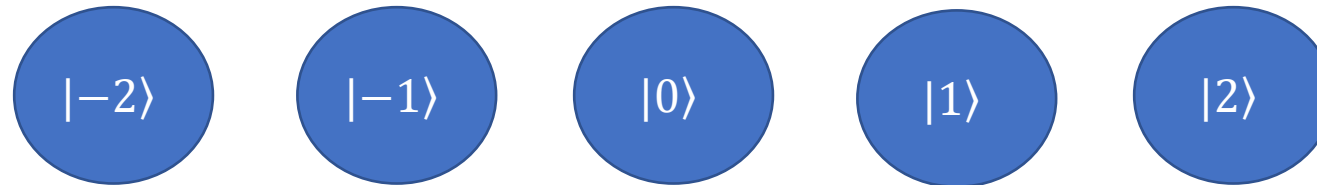


- Asymptotic probability distribution is Gaussian with the position of the peak determined by the bias of the coin.



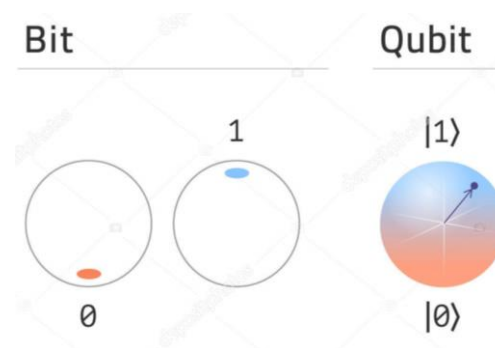
# Some Postulates of Quantum Mechanics

- Hilbert space is a vector space of quantum states (represented by kets- $|\psi\rangle$ ) that is complete.
- Orthonormal basis states representing the position of the particle:  $\{|i\rangle\}$



- Orthonormality of the basis quantum states  $\langle i|j\rangle = \delta_{ij}$
- Uniquely quantum - Superposition of states is also a state in the Hilbert space.

$\{|0\rangle, |1\rangle\}$



$\{\alpha|0\rangle + \beta|1\rangle\}$

# Some Postulates of Quantum Mechanics

- Observable of a physical system is described by an operator that acts on the kets.

$$\hat{p}_x|\psi\rangle = -i\hbar\frac{\partial}{\partial x}|\psi\rangle$$

- The only possible result of the measurement of an observable  $A$  is one of the eigenvalues of the corresponding operator

$$A|\psi_a\rangle = a|\psi_a\rangle$$

# Some Postulates of Quantum Mechanics

- Immediately after the measurement of an observable  $A$  has yielded a value  $a_n$ , the state of the system is the normalized eigenstate  $|\psi_{a_n}\rangle$ .
- The time evolution of a quantum system is given by some unitary operator  $U$  that preserves the normalization of the associated ket.

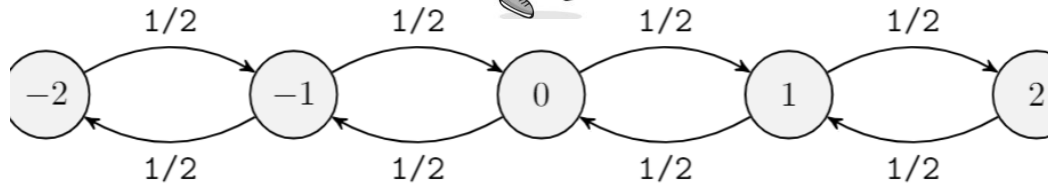
$$|\psi(t)\rangle = U(t, t_0)|\psi(t_0)\rangle$$

i.e,

$$\sum_i |\langle i|\psi(t_0)\rangle|^2 = 1 = \sum_i |\langle i|\psi(t)\rangle|^2$$

# Quantizing the Random Walk

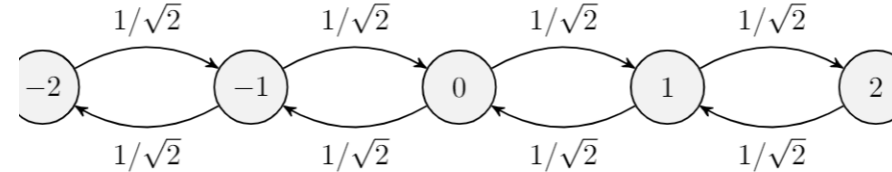
## Classical Random Walk



Matrix elements of  $T$  between Markov States

- Probability distribution after one transition  
 $\mu(t + 1) = \mu(t)T$  , where  $\mu \in \mathbb{R}^N$
- Such that  $\sum_i \mu_i(t) = 1 = \sum_i \mu_i(t + 1)$

## Quantum Walk



Matrix elements between basis quantum states

- Probability distribution after one transition  
 $|\psi(t + 1)\rangle = U\psi(t)$ , where  $\psi \in \mathbb{C}^N$
- Such that ,  
 $\sum_i |\langle i|\psi(t)\rangle|^2 = 1 = \sum_i |\langle i|\psi(t + 1)\rangle|^2$



# Quantizing the Random Walk-Unitary Operator

Coin operator-Direction movement

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$


$$U = (I \otimes H) \cdot S$$

Shift operator that moves the object


$$S = |R\rangle\langle R| \otimes \sum_j |j+1\rangle\langle j| + |L\rangle\langle L| \otimes \sum_j |j\rangle\langle j-1|$$

$$S = S^R + S^L$$

$$U = (I \otimes H)S^R + (I \otimes H)S^L$$

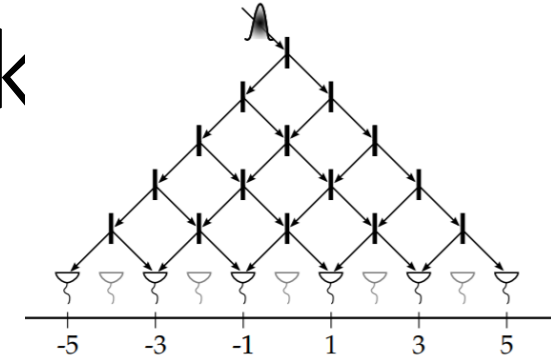


$$T^{(R)} = p \begin{bmatrix} 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 \end{bmatrix}$$



$$T^{(L)} = (1-p) \begin{bmatrix} 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \end{bmatrix}$$

# Mapping the steps of a Quantum Walk



The optical Galton board implementation of quantum walks on the line. The thick lines represent beamsplitters.

- If the object starts in the state

$$\psi(0) = |0R\rangle$$

- First step,  $\psi(1) = U\psi(0)$ ,

Applying  $H \longrightarrow \frac{1}{\sqrt{2}}(|0R\rangle + |0L\rangle)$

$$\psi(1) = \frac{1}{\sqrt{2}}| - 1L\rangle + \frac{1}{\sqrt{2}}|1R\rangle$$

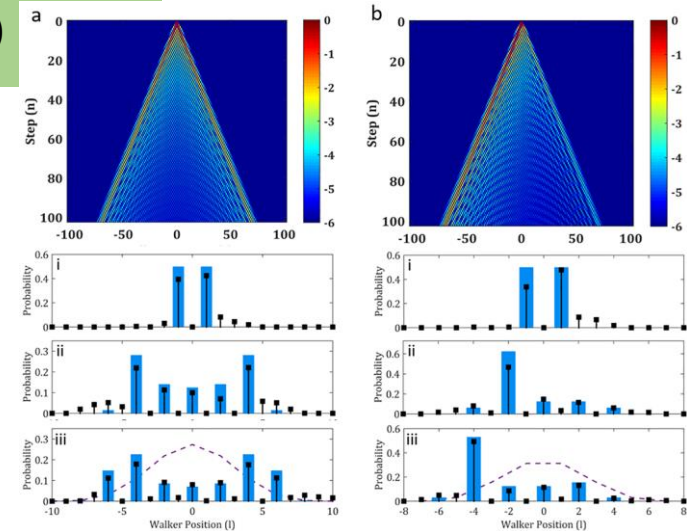
Applying  $S \longrightarrow \frac{1}{\sqrt{2}}(|1R\rangle + | - 1L\rangle)$

- Second step,  $\psi(2) = U\psi(1)$ ,

$$\psi(2) = \frac{1}{2}(| - 2L\rangle + |0L\rangle + |0R\rangle - |2R\rangle)$$

- Third step,  $\psi(3) = U\psi(2)$ ,

$$\psi(3) = \frac{1}{\sqrt{8}}(| - 3L\rangle + | - 1R\rangle + 2|1R\rangle + |1L\rangle + |3R\rangle)$$



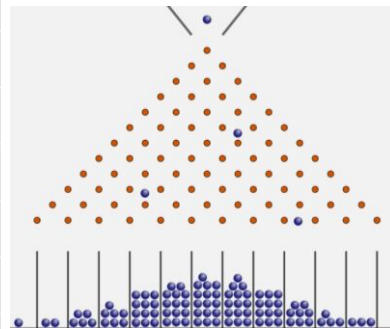
# Random walk vs Quantum Walk

## Random Walk

- The probability of being at position  $i$  after  $T$  steps of the classical random walk on the line starting in 0.



$T \backslash i$	-5	-4	-3	-2	-1	0	1	2	3	4	5
0						1					
1				$\frac{1}{2}$		$\frac{1}{2}$					
2			$\frac{1}{4}$		$\frac{1}{2}$		$\frac{1}{4}$				
3		$\frac{1}{8}$		$\frac{3}{8}$		$\frac{3}{8}$		$\frac{1}{8}$			
4		$\frac{1}{16}$	$\frac{1}{4}$		$\frac{3}{8}$		$\frac{1}{4}$		$\frac{1}{16}$		
5	$\frac{1}{32}$		$\frac{5}{32}$		$\frac{5}{16}$		$\frac{5}{16}$		$\frac{5}{32}$		$\frac{1}{32}$

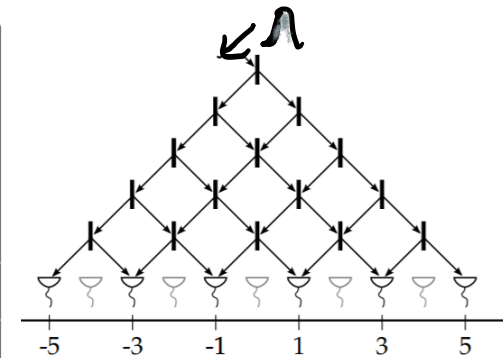


## Quantum Walk

- The probability of being found at position  $i$  after  $T$  steps of the quantum random walk on the line, with the initial state



$T \backslash i$	-5	-4	-3	-2	-1	0	1	2	3	4	5
0						1					
1				$\frac{1}{2}$		$\frac{1}{2}$					
2			$\frac{1}{4}$		$\frac{1}{2}$		$\frac{1}{4}$				
3		$\frac{1}{8}$		$\frac{5}{8}$		$\frac{1}{8}$		$\frac{1}{8}$		$\frac{1}{8}$	
4		$\frac{1}{16}$	$\frac{5}{8}$		$\frac{1}{8}$		$\frac{1}{8}$		$\frac{1}{16}$		
5	$\frac{1}{32}$		$\frac{17}{32}$		$\frac{1}{8}$		$\frac{1}{8}$		$\frac{5}{32}$		$\frac{1}{32}$

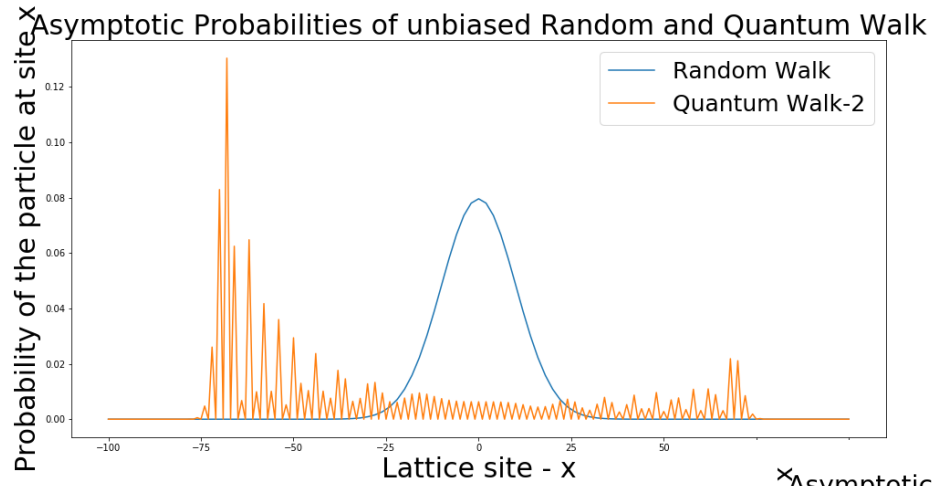


# Quantum Walk

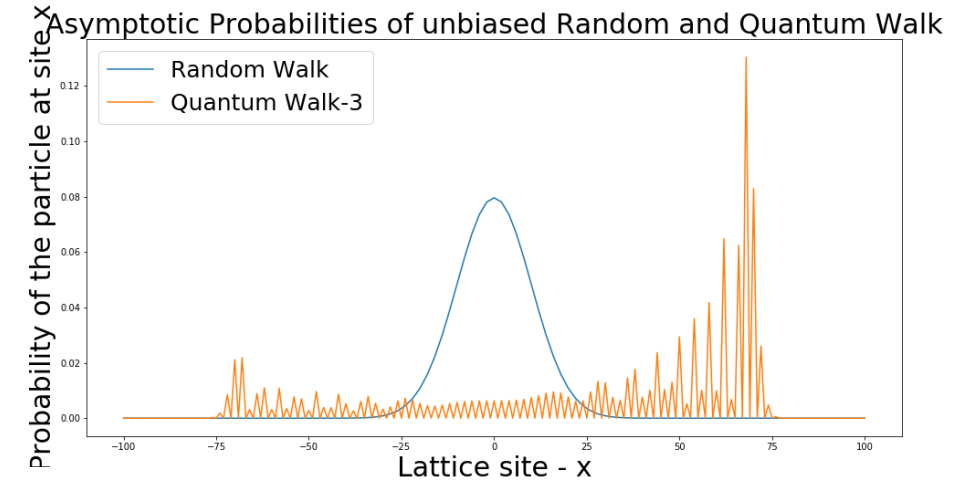
- Asymptotic probability distribution of single particle quantum walk
- Eigenstates of the Unitary operator -  $U = (I \otimes H) \cdot S$
- Eigenvalues of a Unitary operator are of the form  $U|\psi_\lambda\rangle = e^{i\lambda}|\psi_\lambda\rangle$
- The unitary operator has multiple eigenvalues with absolute value 1, therefore there are multiple asymptotic probability distributions

# Quantum Walk

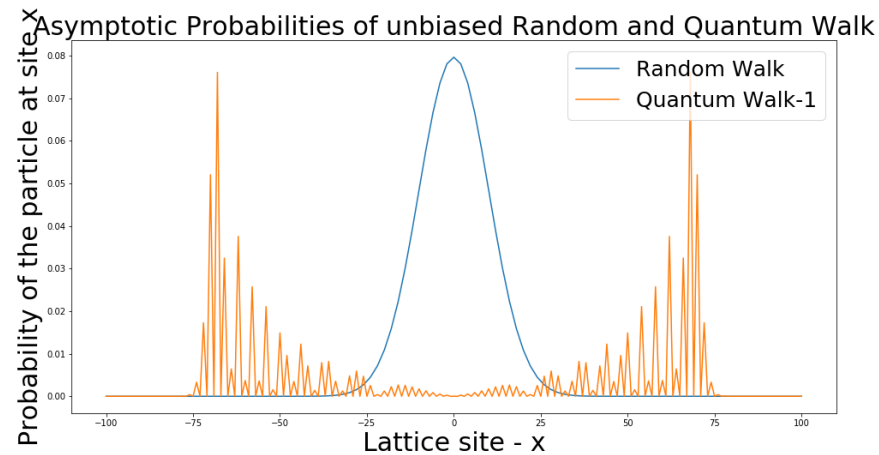
- Direction of the initial state determines which asymptotic state is attained.



$$\psi(0) = |0L\rangle$$



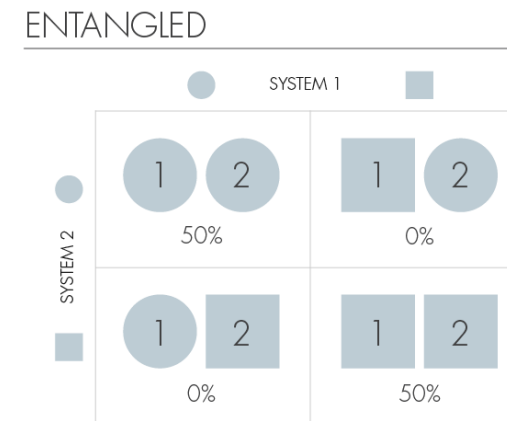
$$\psi(0) = |0R\rangle$$



$$\psi(0) = \frac{1}{\sqrt{2}} (|0R\rangle + i|0L\rangle)$$

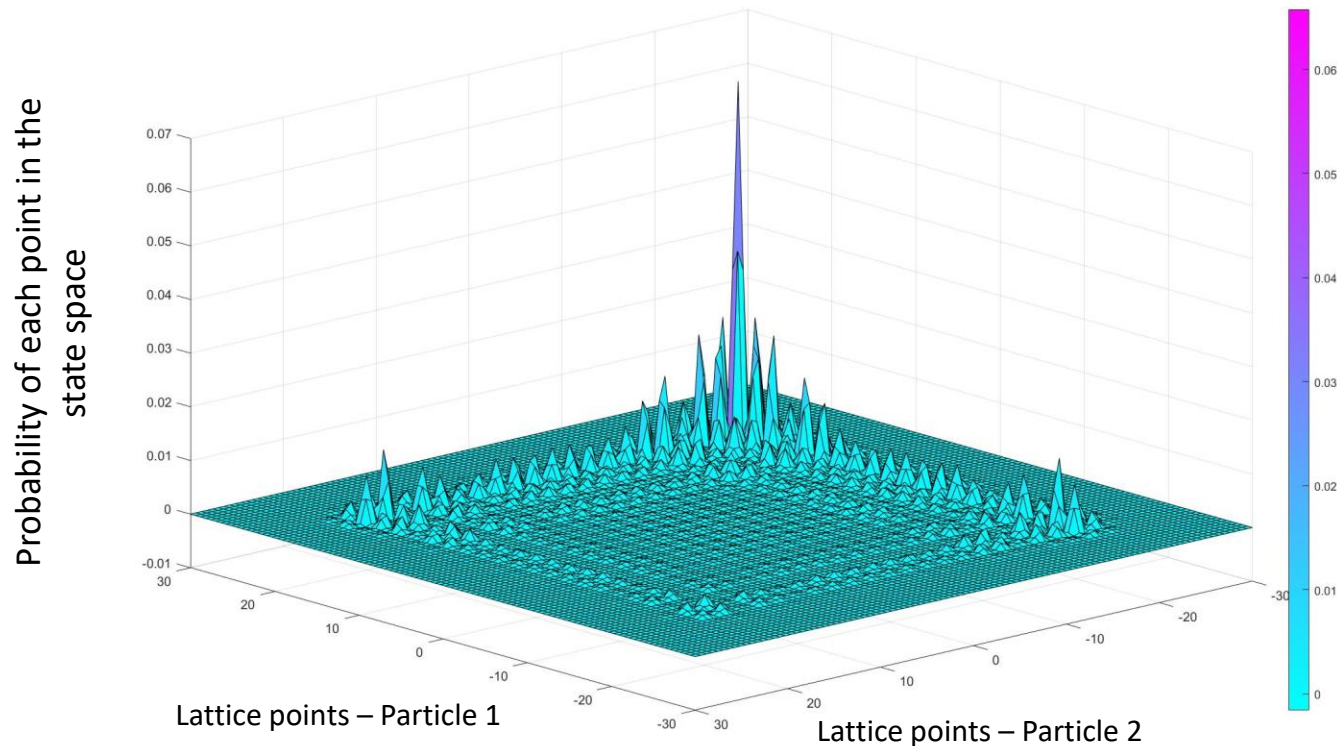
# Two-particle Quantum Walk – non-interacting

- The unitary operator for two non-interacting Quantum walkers is
- $U = U_1 \otimes U_2$  ,  $U_1 = (I \otimes H_1) \cdot S_1, U_2 = (I \otimes H_2) \cdot S_2$
- Two-particle quantum walk allows us to study the effects of indistinguishability and entanglement (another uniquely quantum phenomena (allegedly)).



# Two-particle Quantum Walk – non-interacting

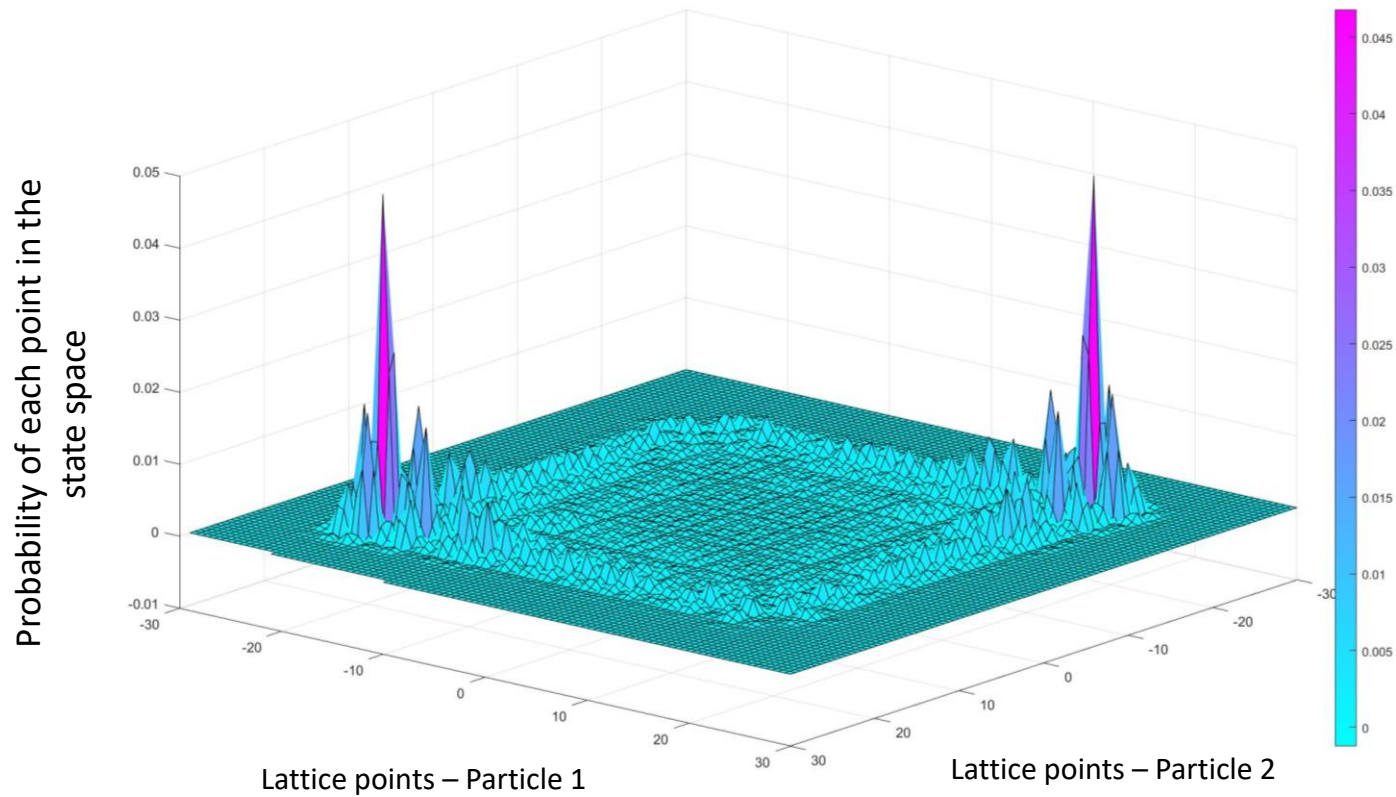
- Numerically obtained probability distribution for two non-interacting Quantum walkers starting in pure state  $\psi(0) = |0L\rangle_1|0R\rangle_2$ .



The asymptotic probability distribution is product of the single particle probability distribution with one walker drifting to left and the other to right

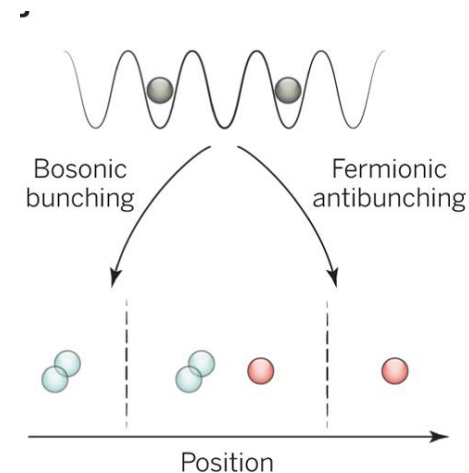
# Two-particle Quantum Walk

- Numerically obtained probability distribution for two non-interacting quantum walker starting in the entangled antisymmetric fermionic state



$$\psi(0) = \frac{1}{\sqrt{2}} (|0L\rangle_1 |0R\rangle_2 - |0R\rangle_1 |0L\rangle_2)$$

The asymptotic probability distribution is not a simple product of the single particle probability distribution. We see that the fermions move away from each other.

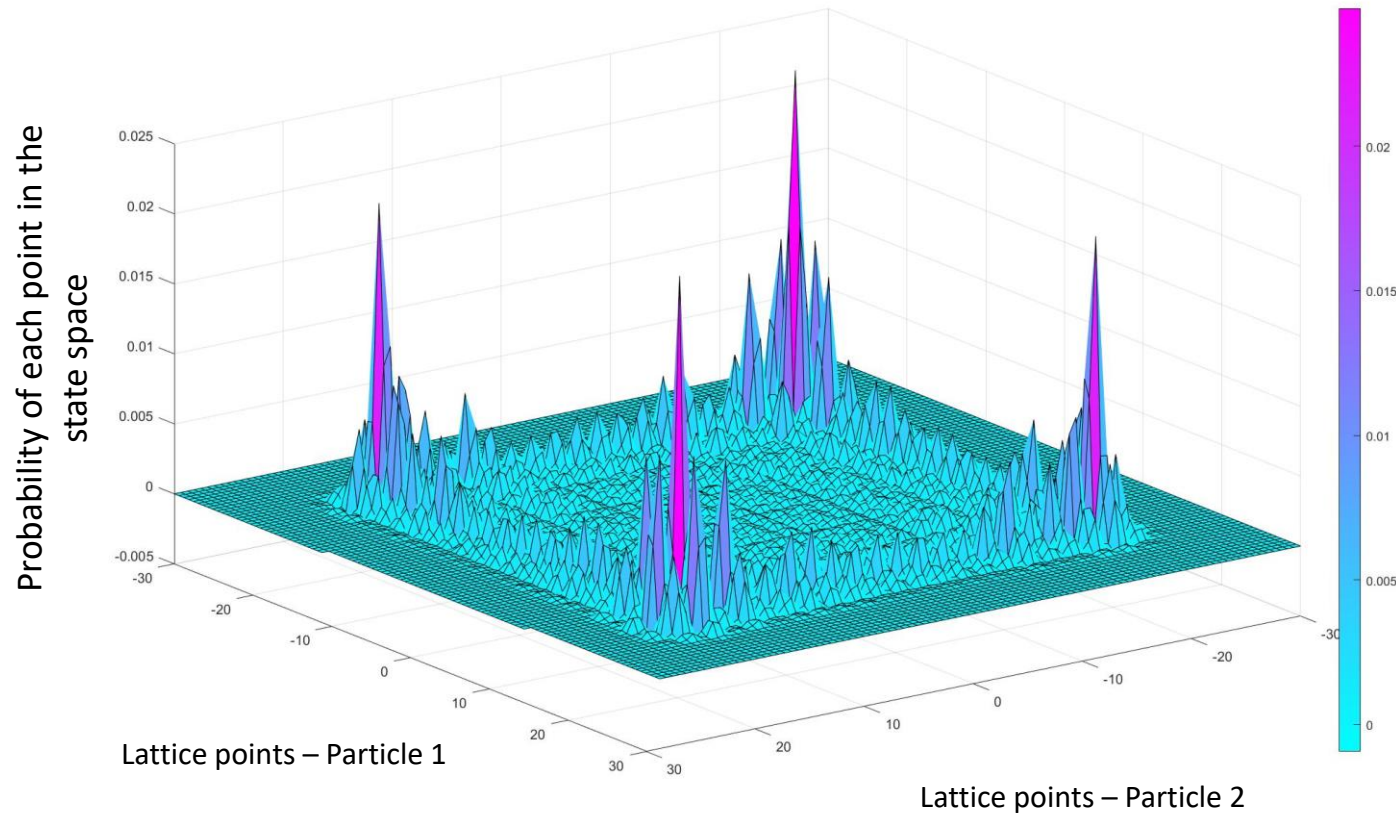




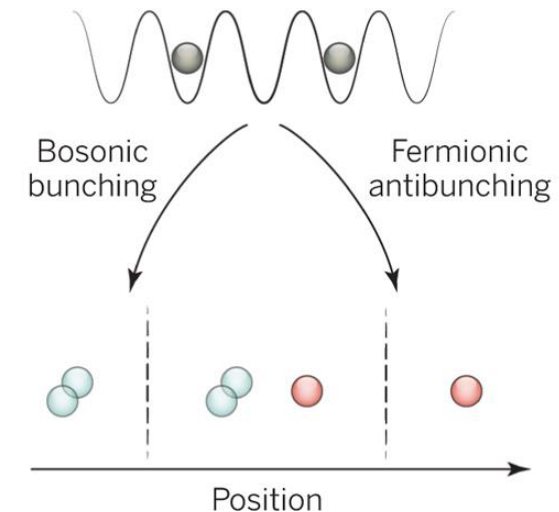
# Two-particle Quantum Walk

- Numerically obtained probability distribution for two non-interacting quantum walker starting in the entangled symmetric boson state

$$\psi(0) = \frac{1}{\sqrt{2}}(|0L\rangle_1|0R\rangle_2 + |0R\rangle_1|0L\rangle_2)$$



The asymptotic probability distribution is not a simple product of the single particle probability distribution. We see that the bosons move towards each other.



# Future work

- Quantum Markov Chain construction of the quantum walk.
- Compare the hitting time for classical and quantum random walkers on a graph using this construction
- Effect of the coin operator on the hitting time.
- Compare and calculate correlations and entropy rate of multiparticle quantum walks especially examining the role of entanglement.
- Entropy rate when the two coins  $H_1$  and  $H_2$  are entangled

Thank You

# Mixed state vs Superposition of States

## Mixed State

- Density Matrix of a mixed state.

$$\rho = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$$

$$\rho = \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$



## Superposition of states

- Density matrix of a superposed state.

$$\rho = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \frac{1}{\sqrt{2}}(\langle 0| + \langle 1|)$$

$$\rho = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$$



# Two Particle-Random Walk

- The state space for a 2D Random Walk:  $\{[i, j]\} \quad i, j \in [-N, N]$

$$T_1 = T_2 = \begin{bmatrix} 0 & 1 & 0 & 0 & 0 \\ 1-p & 0 & p & 0 & 0 \\ 0 & 1-p & 0 & p & 0 \\ 0 & 0 & 1-p & 0 & p \\ 0 & 0 & 0 & 1 & 0 \end{bmatrix}$$

- Transition matrix for a 2D Random walk.

$$T = T_1 \otimes T_2$$

- T is then used to obtain the asymptotic probability distribution  $\pi$  where  $\pi \in \mathbb{R}^{N \times N}$

# Quantum Walk

- Asymptotic probability distribution of single particle quantum walk
- Eigenstates of the Unitary operator

- Eigenvalues of a Unitary operator are of the form

$$U|\psi_\lambda\rangle = e^{i\lambda}|\psi_\lambda\rangle$$

- Multiple eigenstates as there are multiple eigenvalues with absolute value 1

$$U = (I \otimes H) \cdot S$$

Coin operator has two eigenstates whose eigenvalues have absolute value 1.

- Diagonalizing  $S$  involves transforming to the momentum basis.
- Shift operator has two eigenstates whose eigenvalues have absolute value 1.