

# Searching the minimal Quantum $\epsilon$ - Machine

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The presentation is organized as follows:

- $\epsilon$  Machine and formalism of Quantum  $\epsilon$  Machine
- Examples of Quantum  $\epsilon$  Machine
- Discussion and future scope of work

# $\epsilon$ Machine

- Causal states:  $s_i, i = 1, \dots, n$
- Output alphabet:  $\mathcal{A}$
- Transition matrices:  $T_{ij}^a, a \in \mathcal{A}$
- Property : Unifilarity

# Formalism of Quantum $\epsilon$ Machine

- Causal states:  $|\psi_i\rangle$ ,  $i = 1, \dots, n$  (Normalized)
- Operators:  $\mathcal{T}^a$ ,  $a \in \mathcal{A}$
- Unifilarity:  $\mathcal{T}^a|\psi_i\rangle = U_{ij}^a|\psi_j\rangle \sim |\psi_j\rangle$

Probability: Suppose we are currently in state  $\psi_i$ , the probability of observing 01 is given by:

$$\text{Prob} = \sum_j |\langle\psi_i|\mathcal{T}^0\mathcal{T}^1|e_j\rangle|^2 \quad (1)$$

where  $e_j$  are basis vectors. We can rewrite the probability as:

$$\text{Tr}(|\psi_i\rangle\langle\psi_i|\mathcal{T}^0\mathcal{T}^1(\mathcal{T}^1)^\dagger(\mathcal{T}^0)^\dagger) \quad (2)$$

In general if we are in  $\rho$  the probability of observing  $a_1 \dots a_n$  is:

$$\text{Tr}(\rho\mathcal{T}^{a_1} \dots \mathcal{T}^{a_n}(\mathcal{T}^{a_n})^\dagger \dots (\mathcal{T}^{a_1})^\dagger) \quad (3)$$

# Properties of Operators and States

- For any state  $\rho$ , the total probability of observing all outputs is 1:

$$1 = \sum_a \text{Tr}(\rho \mathcal{T}^a (\mathcal{T}^a)^\dagger) \quad (4)$$

This is true for any states so we must ask

$$\sum_a \mathcal{T}^a (\mathcal{T}^a)^\dagger = \mathbf{1} \quad (5)$$

# Properties of Operators and States

- The transition probabilities are encoded in this way:

$$\text{Prob}(i \rightarrow a, j) = \text{Tr}(|i\rangle\langle i| \mathcal{T}^a (\mathcal{T}^a)^\dagger) = |U_{ij}^a|^2 = \text{Pr}(j, a|i) \quad (6)$$

In general we have:

$$\mathcal{T}^a |i\rangle = \sqrt{\text{Pr}(j, a|i)} \exp(i\theta_{ij}^a) |j\rangle \quad (7)$$

- The inner product of two states are given by:

$$\begin{aligned} \langle i|j\rangle &= \langle i|\mathbf{1}|j\rangle = \langle i| \sum_a \mathcal{T}^a (\mathcal{T}^a)^\dagger |j\rangle \\ &= \sum_a \sqrt{\text{Pr}(l, a|j) \text{Pr}(k, a|i)} (i\theta_{jl}^a - i\theta_{ik}^a) \langle k|l\rangle \end{aligned} \quad (8)$$

# Properties of Operators and States

- For each transition  $i \rightarrow j$  with output  $a$  we have a phase parameter  $\theta_{ij}^a$
- We want to find out the parameters that minimize the von Neumann entropy

# Perturbed Coin I

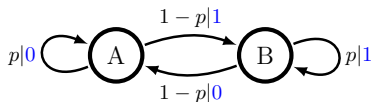


Figure 1: Perturbed Coin

$$\mathcal{T}^0 |A\rangle = \sqrt{p} \exp(i\theta_1) |A\rangle \quad \mathcal{T}^0 |B\rangle = \sqrt{1-p} \exp(i\theta_2) |A\rangle \quad (9)$$

$$\mathcal{T}^1 |B\rangle = \sqrt{p} \exp(i\theta_3) |A\rangle \quad \mathcal{T}^1 |A\rangle = \sqrt{1-p} \exp(i\theta_4) |B\rangle$$

$$\begin{aligned} \langle B|A\rangle &= \sqrt{p}\sqrt{1-p} \left( \exp(i\theta_1 - i\theta_2) + \exp(i\theta_4 - i\theta_3) \right) \\ &= \sqrt{p}\sqrt{1-p} r \exp(i\theta) \end{aligned} \quad (10)$$

where  $r \in [0, 2]$  and  $\theta \in [0, 2\pi]$ .



## Perturbed Coin II

The state we prepare is:

$$\rho = \frac{1}{2}|A\rangle\langle A| + \frac{1}{2}|B\rangle\langle B| \quad (11)$$

The eigenvalues are:

$$\lambda_1 = \frac{1}{2} \left( 1 - r\sqrt{p - p^2} \right), \lambda_2 = \frac{1}{2} \left( 1 + r\sqrt{p - p^2} \right) \quad (12)$$

When  $r = 2$  von Neumann entropy obtains its minimum. They are the same as q-machine:

$$\begin{aligned} |B\rangle &= \sqrt{1-p}|e_0\rangle + \sqrt{p}|e_1\rangle \\ |A\rangle &= \sqrt{p}|e_0\rangle + \sqrt{1-p}|e_1\rangle \end{aligned} \quad (13)$$

where  $|e_0\rangle$  and  $|e_1\rangle$  are two orthogonal states.

# Nemo Process I

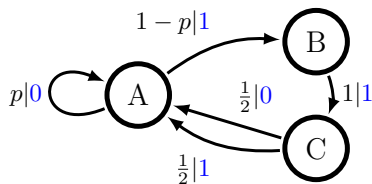


Figure 2: Nemo process

$$\mathcal{T}^0 |A\rangle = \sqrt{p} \exp(i\theta_1) |A\rangle \quad (14)$$

$$\mathcal{T}^0 |B\rangle = 0$$

$$\mathcal{T}^0 |C\rangle = \sqrt{1/2} \exp(i\theta_2) |A\rangle$$

$$\mathcal{T}^1 |A\rangle = \sqrt{1-p} \exp(i\theta_3) |B\rangle$$

$$\mathcal{T}^1 |B\rangle = \exp(i\theta_4) |C\rangle$$

$$\mathcal{T}^1 |C\rangle = \sqrt{1/2} \exp(i\theta_5) |A\rangle$$

$$\begin{aligned}\langle A|B\rangle &= \frac{\sqrt{p(1-p)}}{1+p} \exp(i\theta_1 - \theta_2 - \theta_3 + \theta_5) \\ \langle B|C\rangle &= \frac{\sqrt{p}}{1+p} \exp(i\theta_1 - \theta_2 - \theta_4 + \theta_5) \\ \langle A|C\rangle &= \frac{\sqrt{2p}}{1+p} \exp(i\theta_1 - \theta_2)\end{aligned}\tag{15}$$

We can choose phases such that  $\langle A|B\rangle$  and  $\langle B|C\rangle$  are real numbers:

$$\begin{aligned}\langle A|B\rangle &= \frac{\sqrt{p(1-p)}}{1+p} \\ \langle B|C\rangle &= \frac{\sqrt{p}}{1+p} \\ \langle A|C\rangle &= \frac{\sqrt{2p}}{1+p} \exp(i\theta)\end{aligned}\tag{16}$$

## Nemo Process III

The mixed state we prepare is:

$$\rho = \frac{1}{3-2p}|A\rangle\langle A| + \frac{1-p}{3-2p}|B\rangle\langle B| + \frac{1-p}{3-2p}|C\rangle\langle C| \quad (17)$$

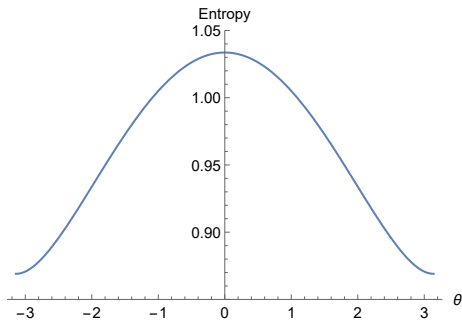


Figure 3: Entropy- $\theta$  for  $p=0.666$  Nemo Process

# Nemo Process IV

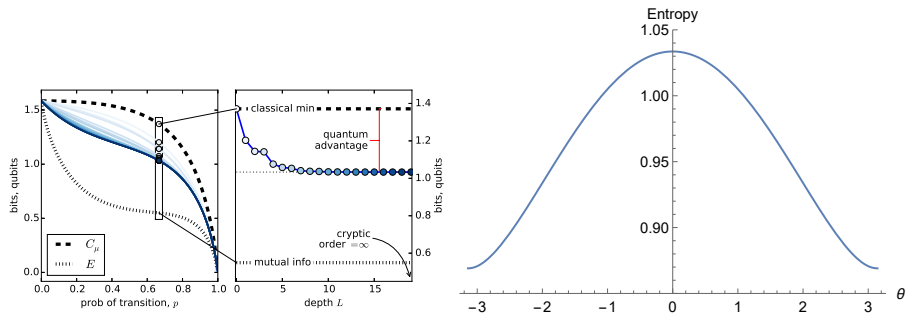


Figure 4: Entropy for  $p=0.666$  Nemo Process. (Mahoney, Aghamohammadi, Crutchfield. 2016)

- For Nemo process we obtain the minimum entropy at

$$\langle A|B \rangle = \frac{\sqrt{p(1-p)}}{1+p} \quad (18)$$

$$\langle B|C \rangle = \frac{\sqrt{p}}{1+p}$$

$$\langle A|C \rangle = -\frac{\sqrt{2p}}{1+p}$$

- When  $p = \frac{\sqrt[3]{9+\sqrt{87}}}{6^{2/3}} - \frac{1}{\sqrt[3]{6(9+\sqrt{87})}} \sim 0.589755$ , we can use only two states to simulate this process.

### 3 States MBW process

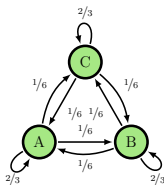


Figure 5: 3 States MBW process

- The inner products are given by:

$$\begin{aligned}\langle A|B\rangle &= r_{AB} \\ \langle B|C\rangle &= r_{BC} \\ \langle A|C\rangle &= r_{AC} \exp(i\theta)\end{aligned}\tag{19}$$

where  $r_{AB}, r_{BC}, r_{AC} \leq \frac{5}{6}$  and  
 $1 + 2 r_{AB} r_{BC} r_{AC} \cos(\theta) \geq r_{AB}^2 + r_{BC}^2 + r_{AC}^2$ .

- The mixed state we prepare is:

$$\rho = \frac{1}{3}|A\rangle\langle A| + \frac{1}{3}|B\rangle\langle B| + \frac{1}{3}|C\rangle\langle C| \quad (20)$$

- The minimum are given at:

$$\begin{aligned} \langle A|B\rangle &= \frac{5}{6} \\ \langle B|C\rangle &= \frac{5}{6} \\ \langle A|C\rangle &= \frac{5}{6} \exp(i\theta) \end{aligned} \quad (21)$$

where  $\cos(\theta) = \frac{117}{125}$

- Three eigenvalues are  $\{0, \frac{1}{18} (9 - 4\sqrt{3}), \frac{1}{18} (9 + 4\sqrt{3})\}$ .
- von Neumann entropy  $S = 0.515109$  bits. For 3D q-machine ( $r_{AB}, r_{BC}, r_{AC} = \frac{5}{6}, \theta = 0$ )  $S_{q3} = 0.61$  bits and for 2D q-machine ( $r_{AB}, r_{BC}, r_{AC} = \frac{1}{2}, \theta = \pi$ )  $S_{q2} = 1$  bits. (Loomis, Crutchfield 2018)



- For larger  $\epsilon$ -Machine, how to find out the parameters that minimize von Neumann entropy
- By using mixed states to represent causal states can we do better?

THANK YOU!